

Astrophysical insights from the compact-object mass distribution inferred with gravitational waves

Sylvia Biscoveanu
NHFP Symposium 2024

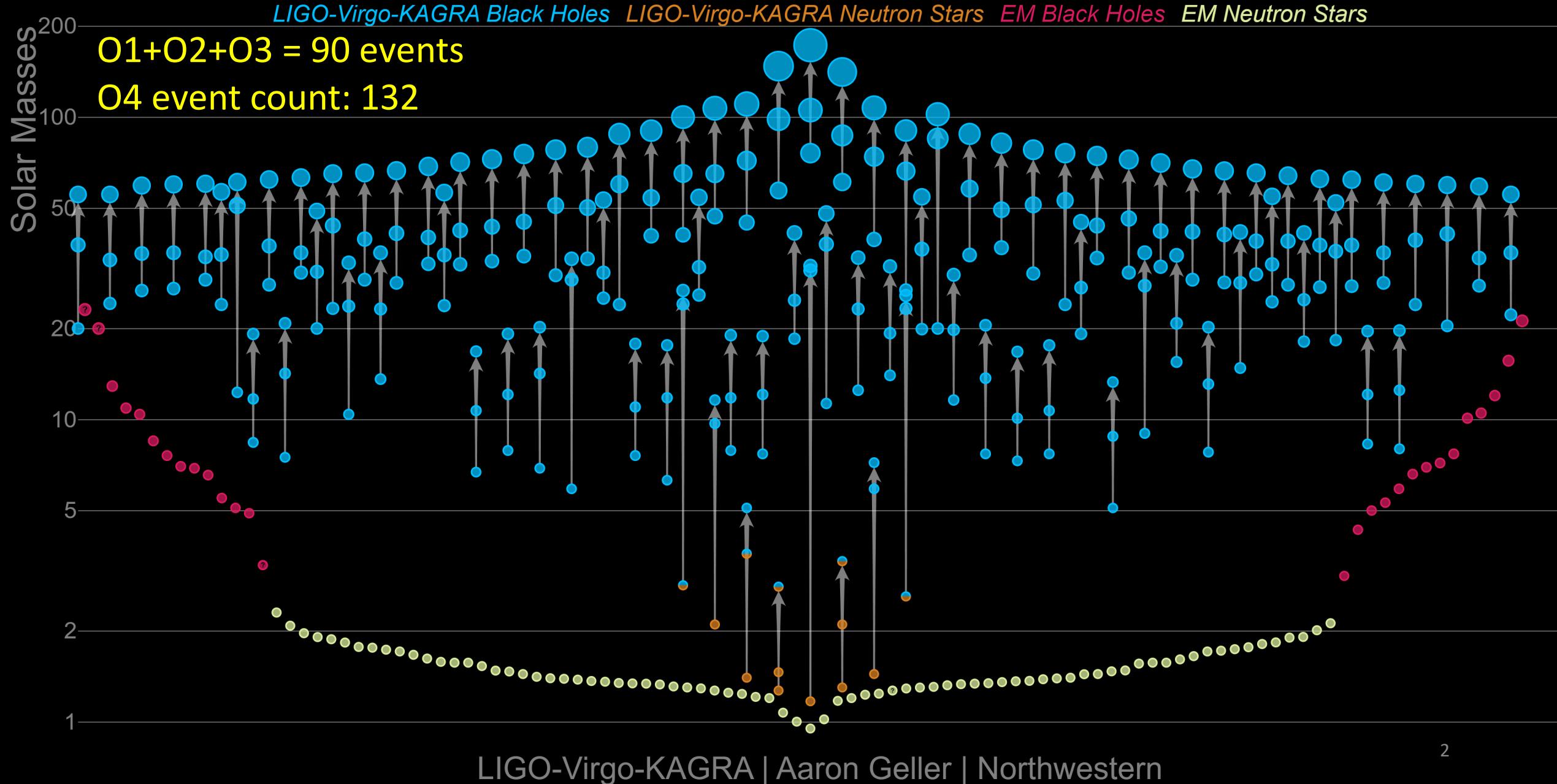


 @sylvia_bisco
 sbisco@northwestern.edu



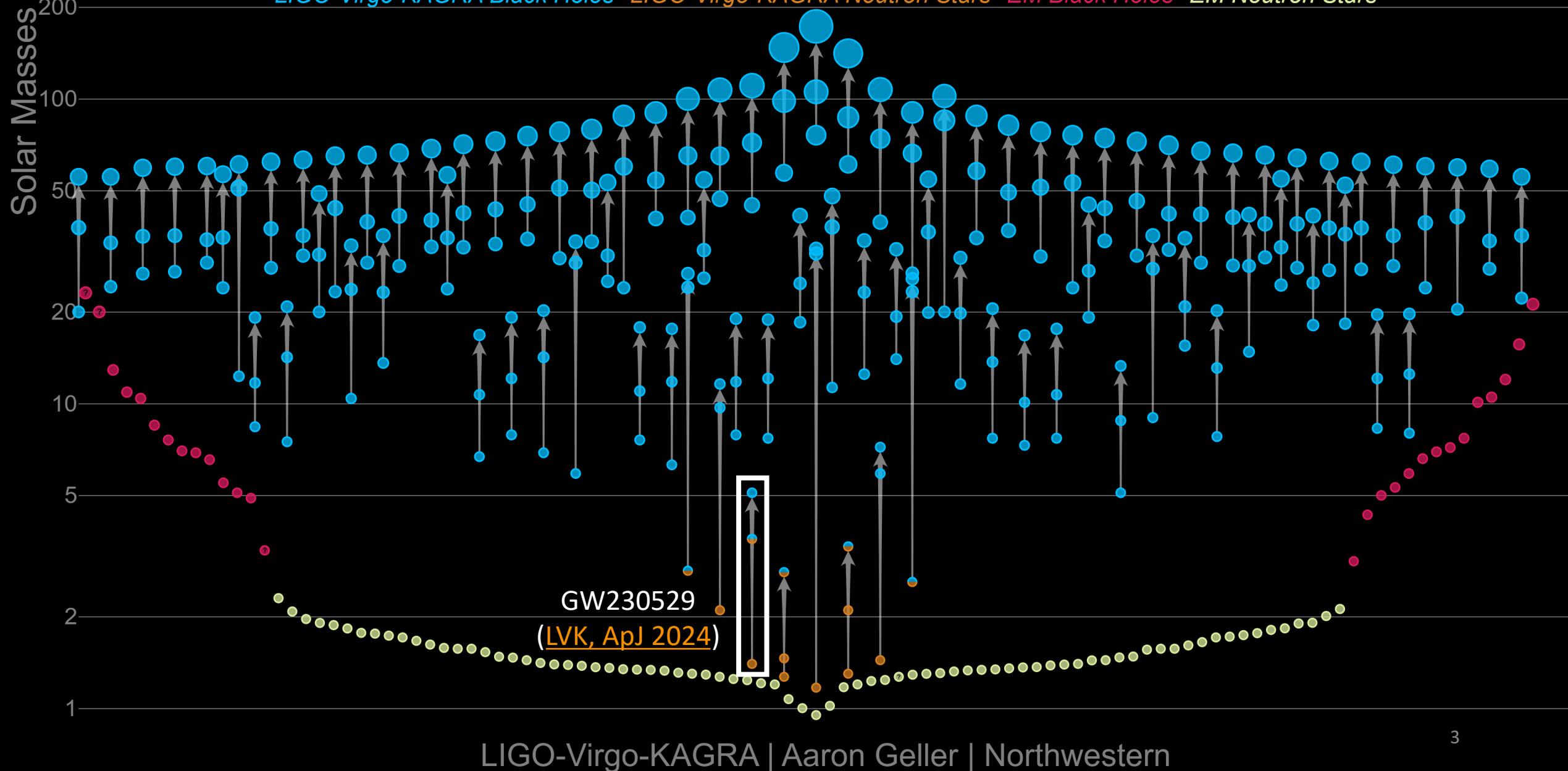
NASA Hubble
Fellowship Program

Masses in the Stellar Graveyard

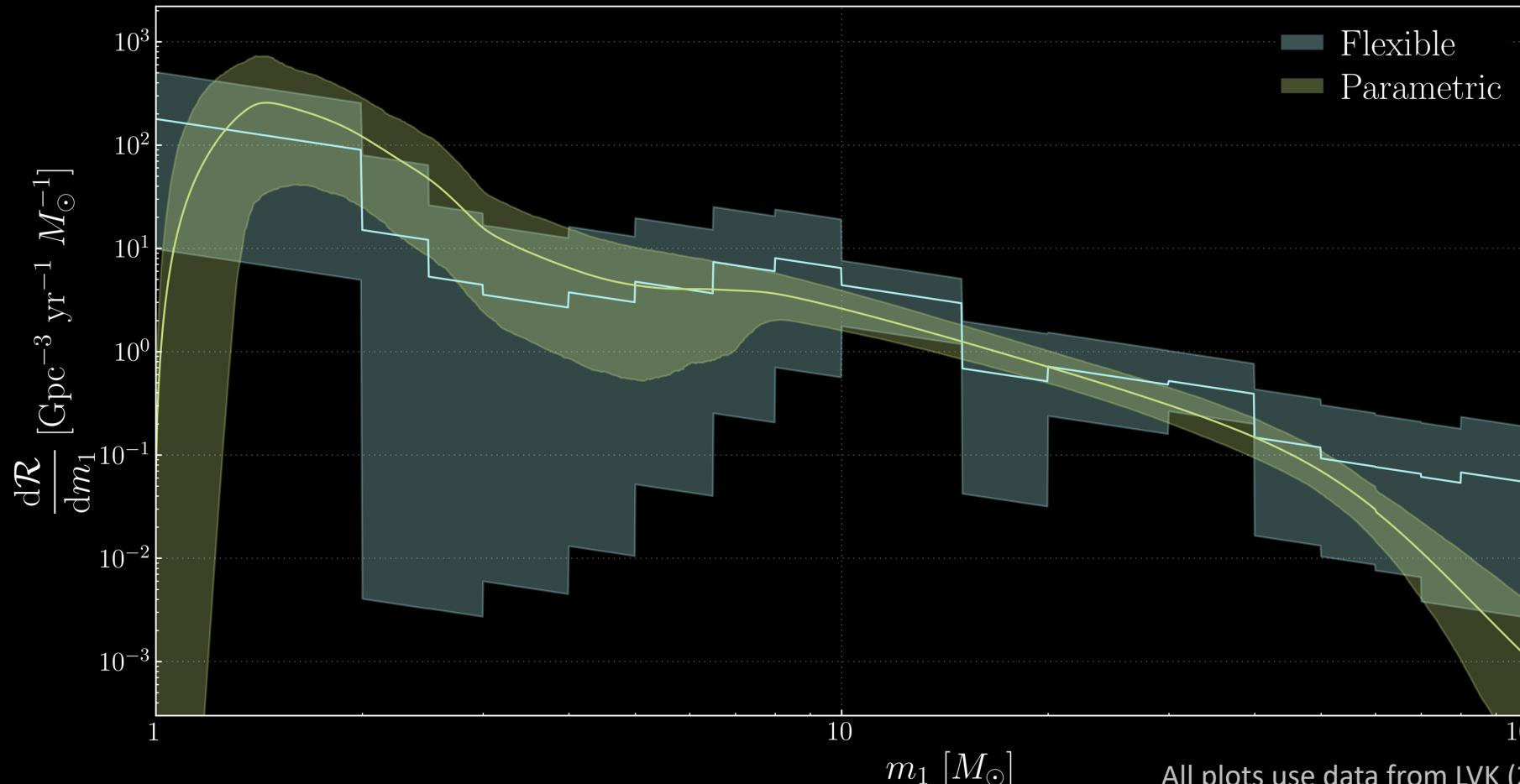


Masses in the Stellar Graveyard

LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars EM Black Holes EM Neutron Stars

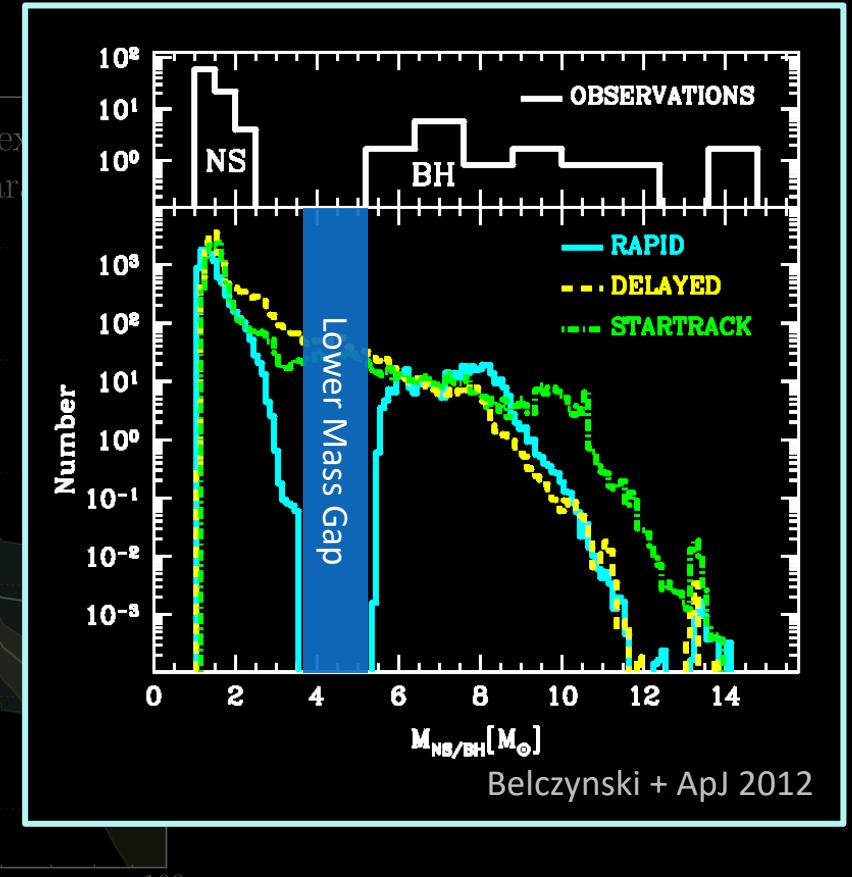
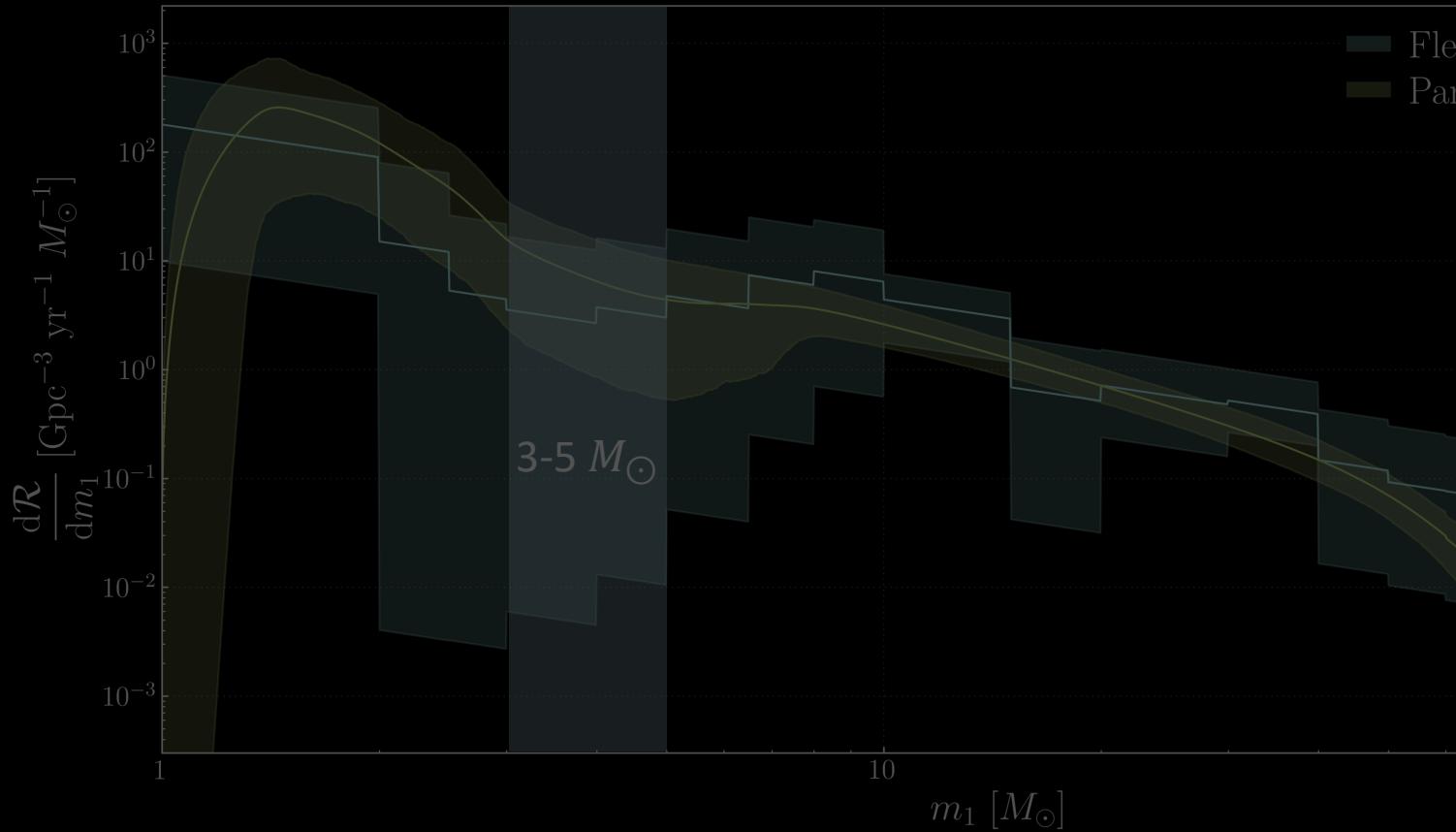


The compact object mass distribution provides key astrophysical insights:

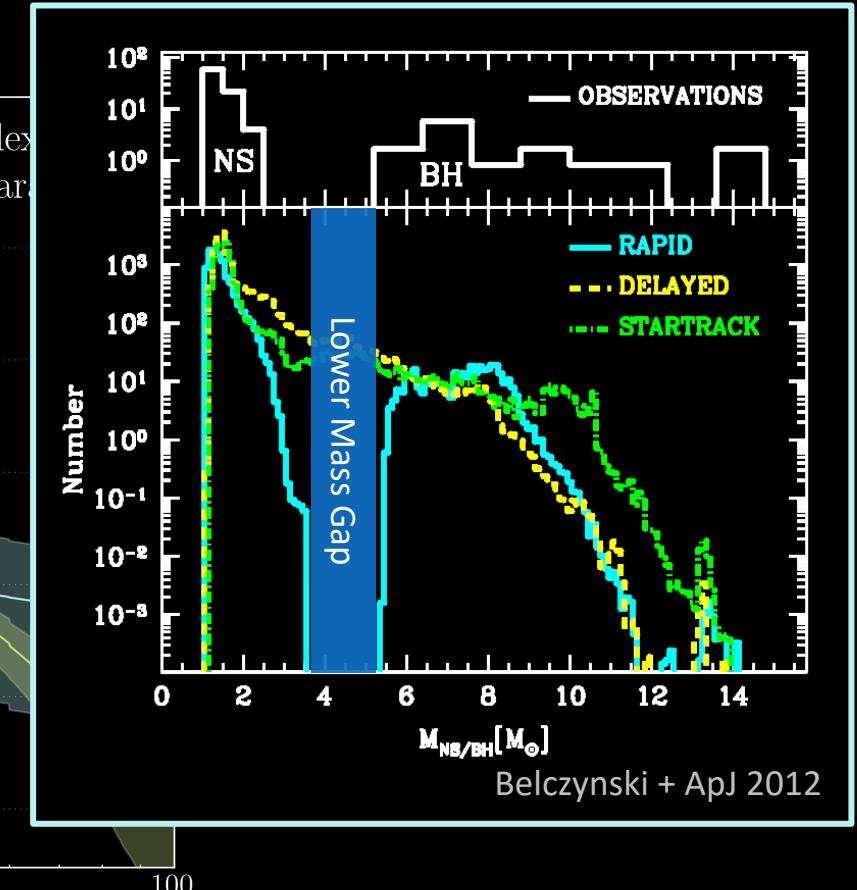
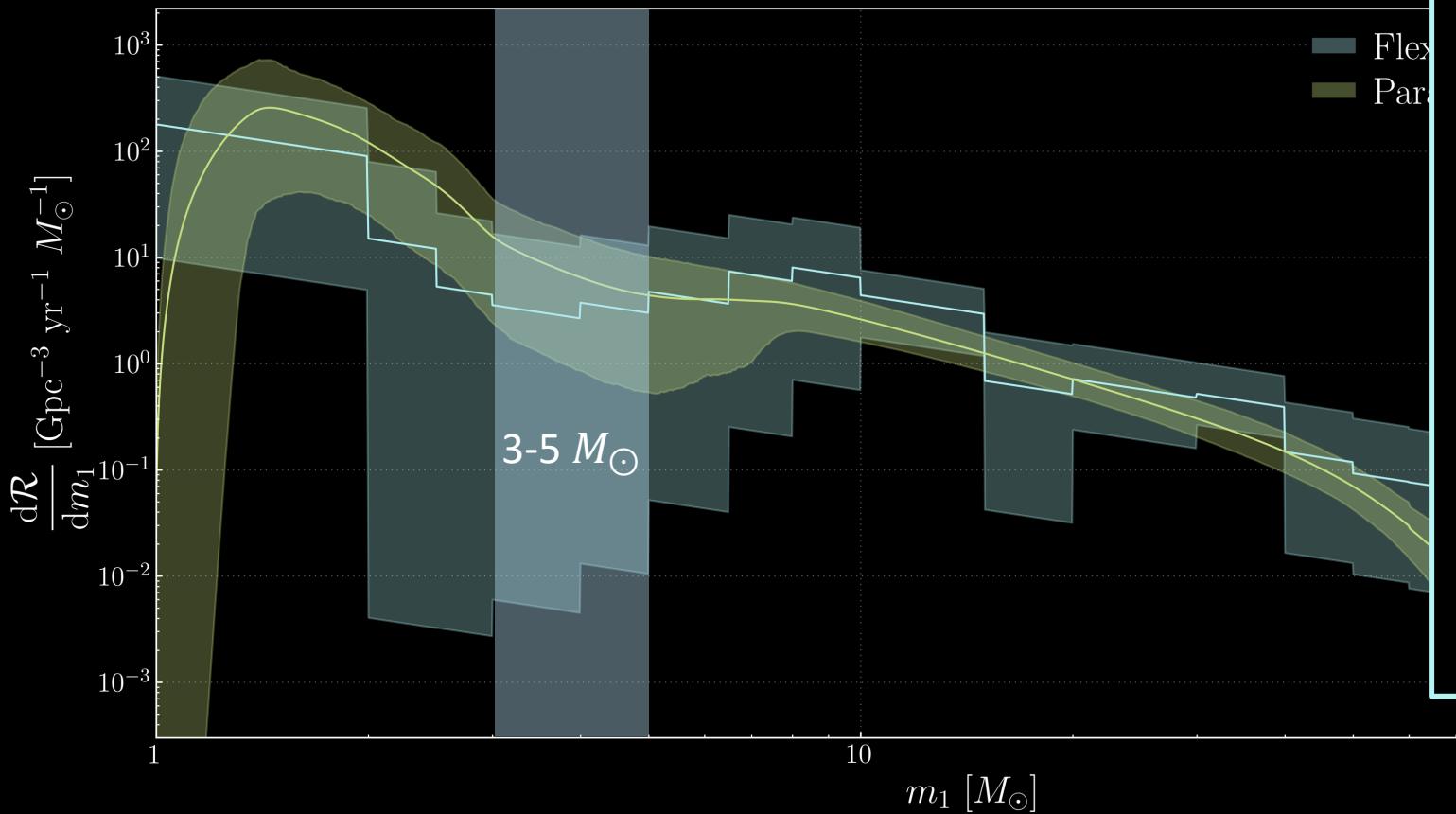


All plots use data from LVK (2024)
unless otherwise noted

Evidence for a gap between the most massive neutron stars and least massive black holes?

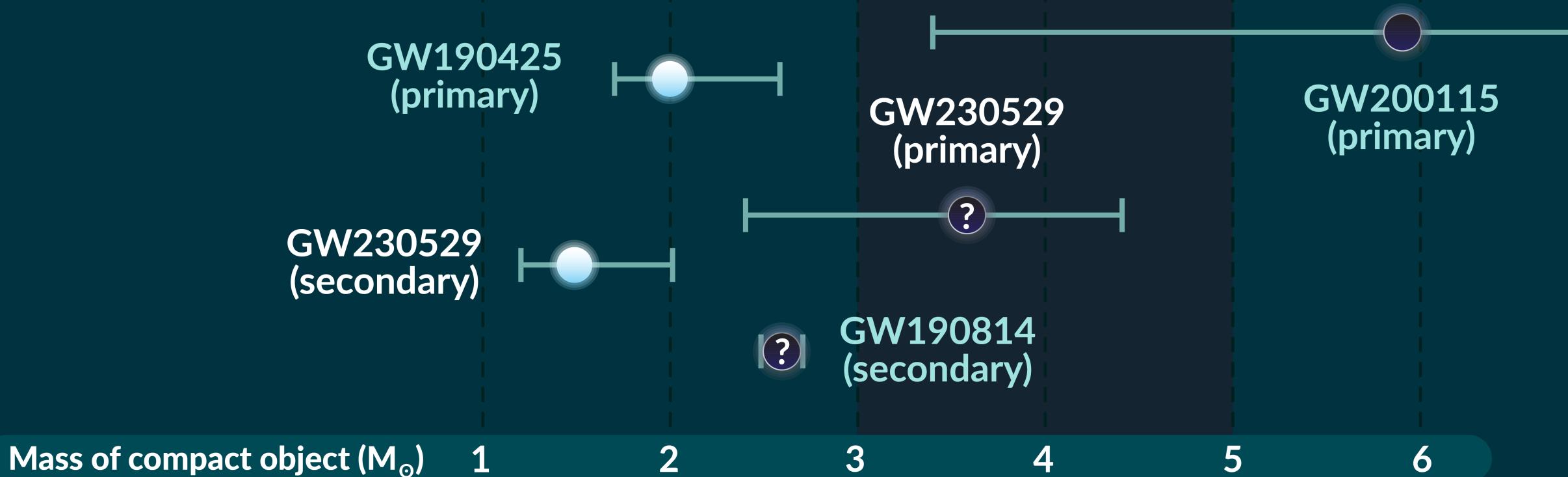


Evidence for a gap between the most massive neutron stars and least massive black holes?



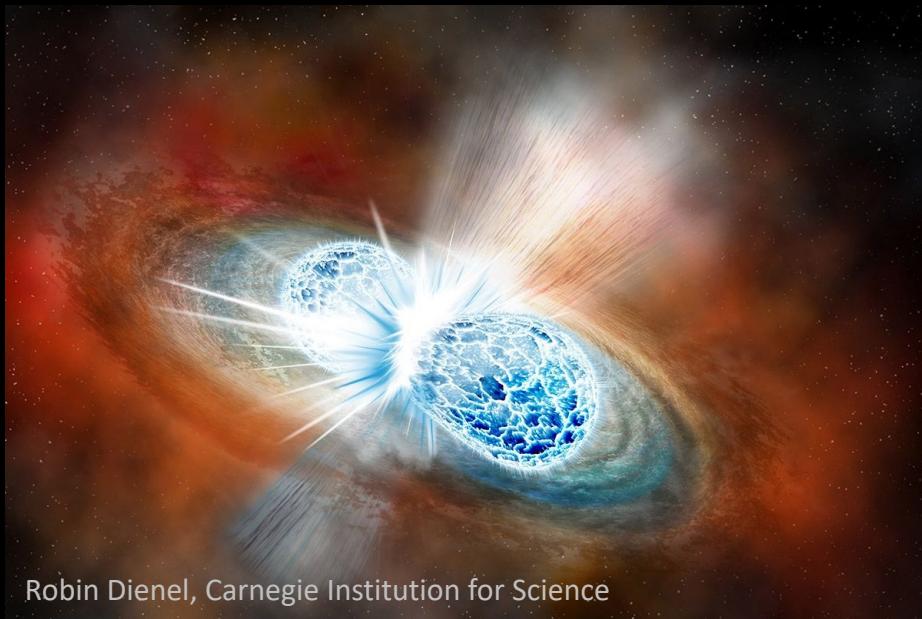
FILLING THE MASS

with observations of compact binaries from gravitational waves



Includes components of compact binary mergers detected with a False Alarm Rate (FAR) of less than 0.25 per year

GW230529: the merger of a neutron star with a lower-mass-gap object



Robin Dienel, Carnegie Institution for Science

Heavy binary neutron
star merger

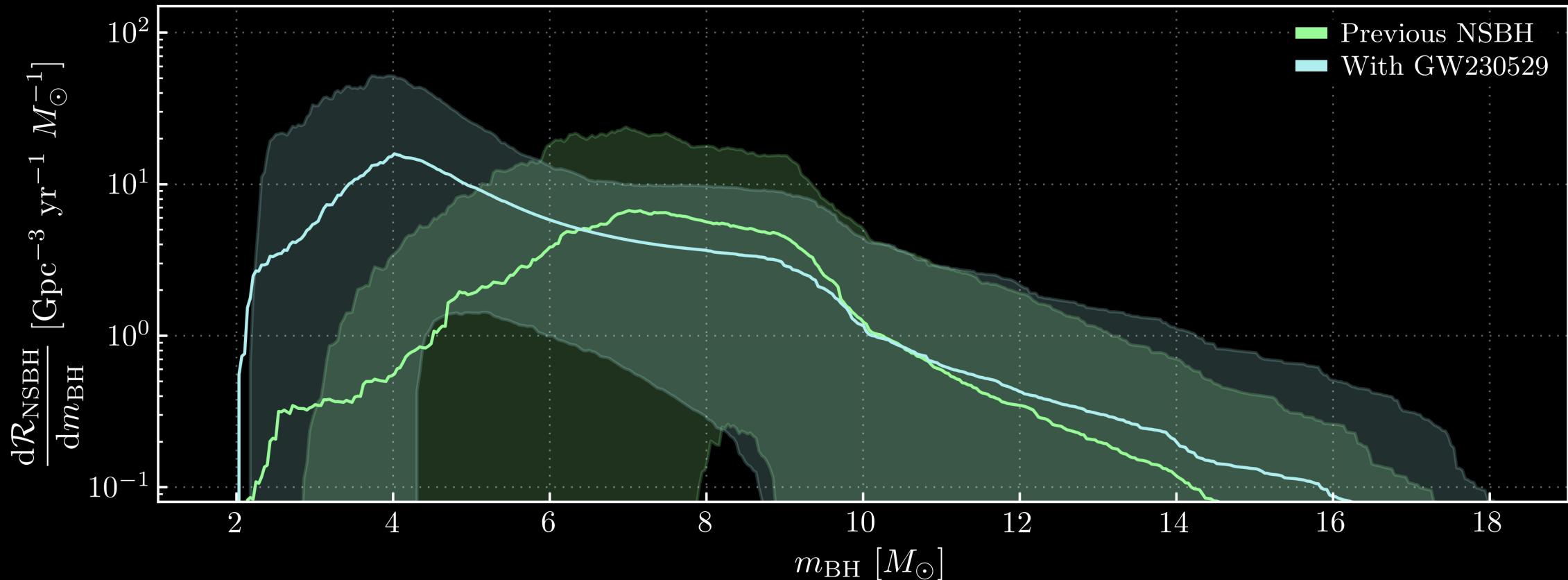
OR



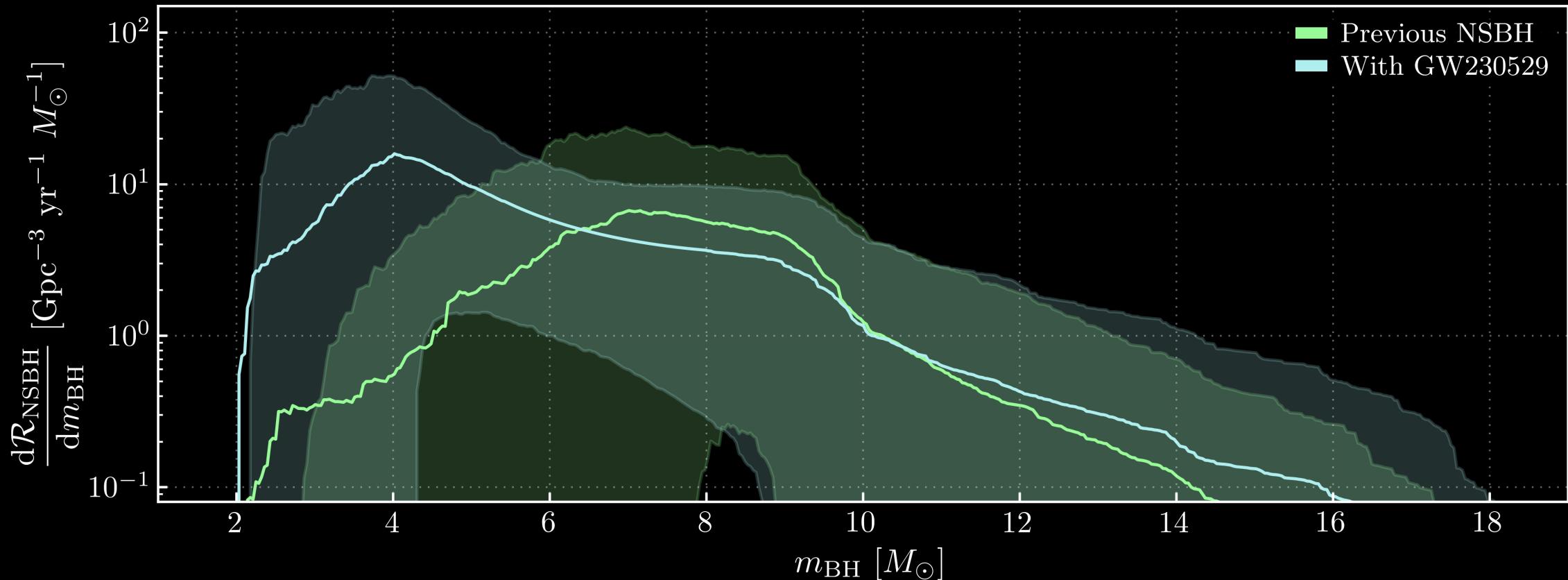
Carl Knox, OzGrav/Swinburne

Neutron star merging
with low-mass black hole

Including GW230529 decreases the minimum mass of black holes in the NSBH population



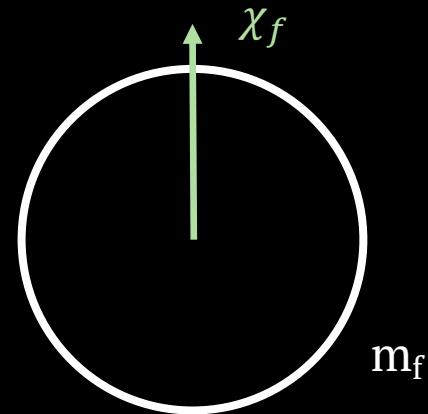
Including GW230529 decreases the minimum mass of black holes in the NSBH population



Probability of the existence of a mass gap (in NSBH) decreases from 98.6% to 7.2%

Merger outcome depends on binary properties

Black hole swallows
neutron star



Rigid neutron stars,
unequal mass ratios, small
black hole spins

Neutron star tidally
disrupted outside ISCO

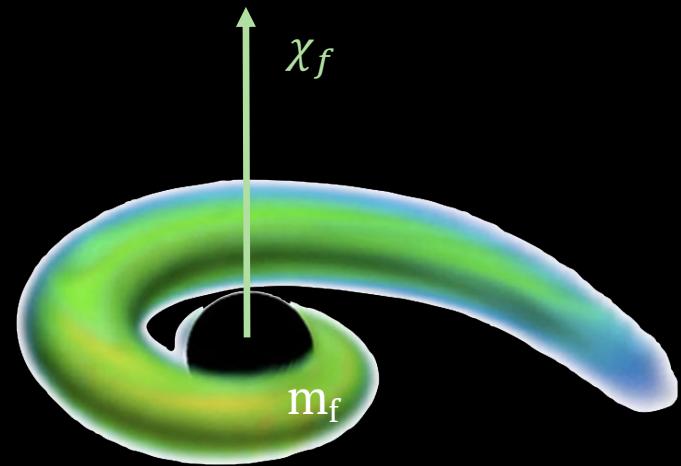
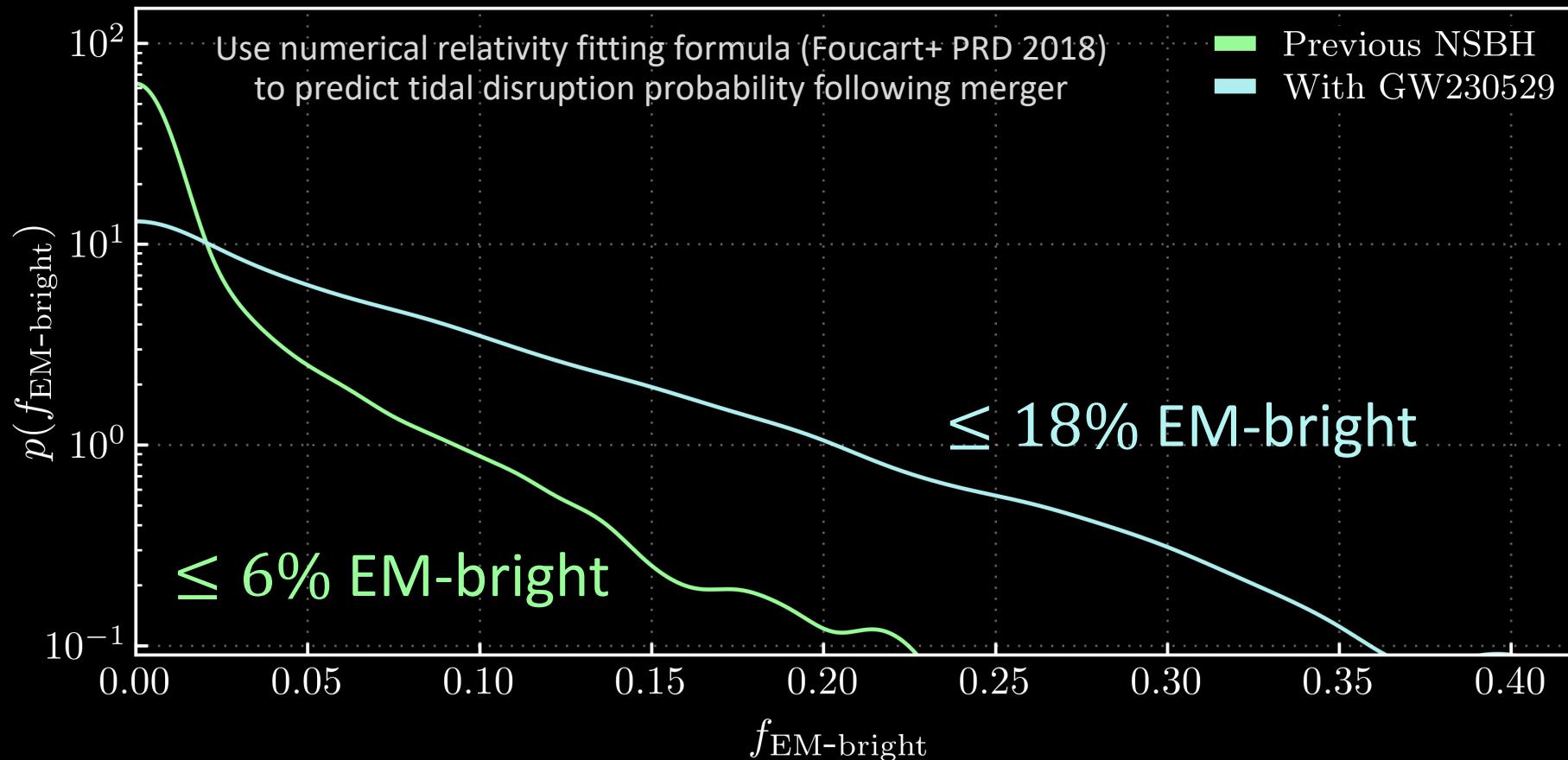


Image: Foucart FrASS 2020

Deformable neutron stars, equal
mass ratios, large prograde spins

The detection of GW230529 means that NSBH mergers are more likely to be multimessenger sources



GW230529 Summary

Likely the merger of a neutron star with a mass-gap black hole

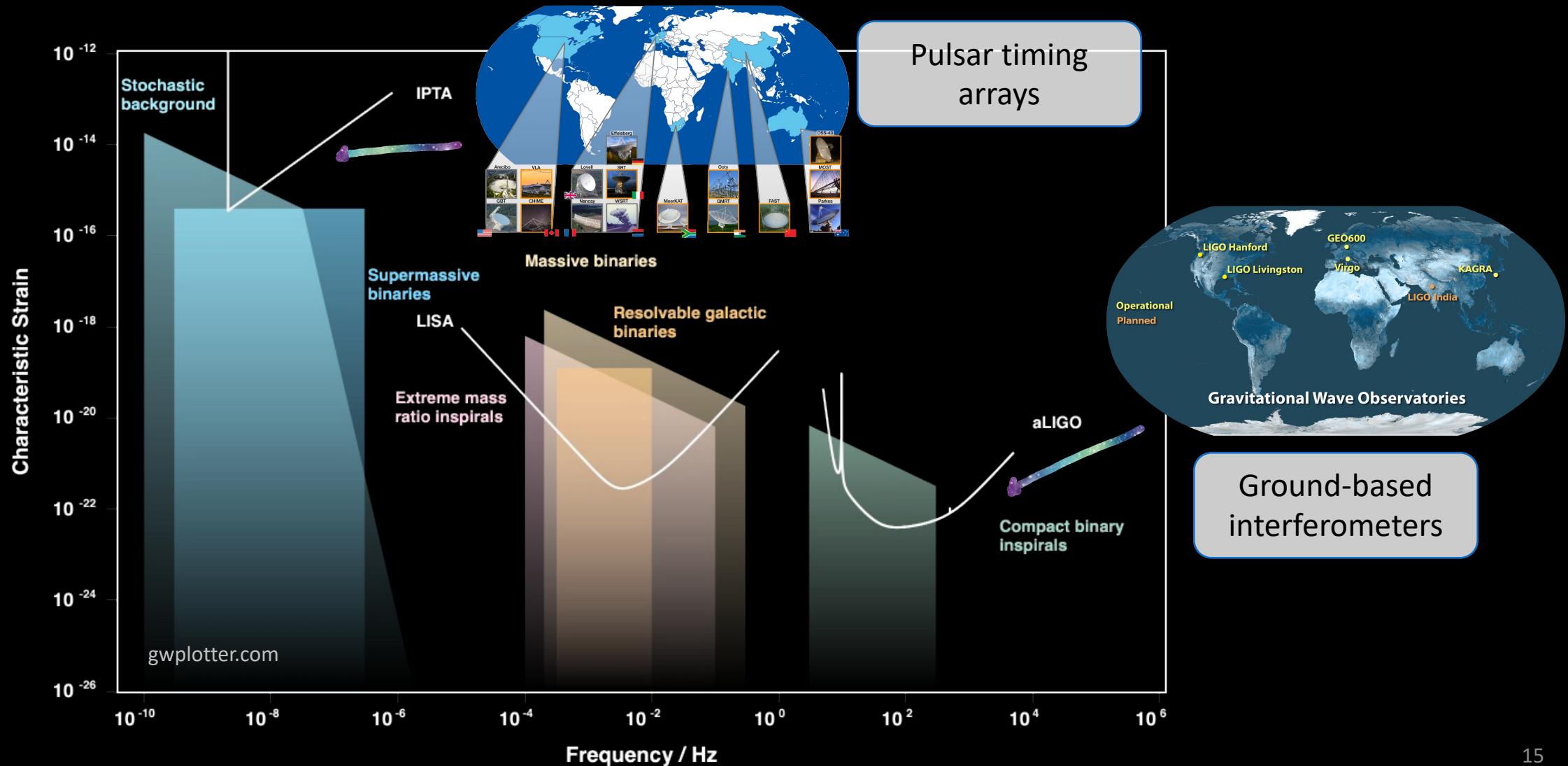
Challenges compact object formation paradigms

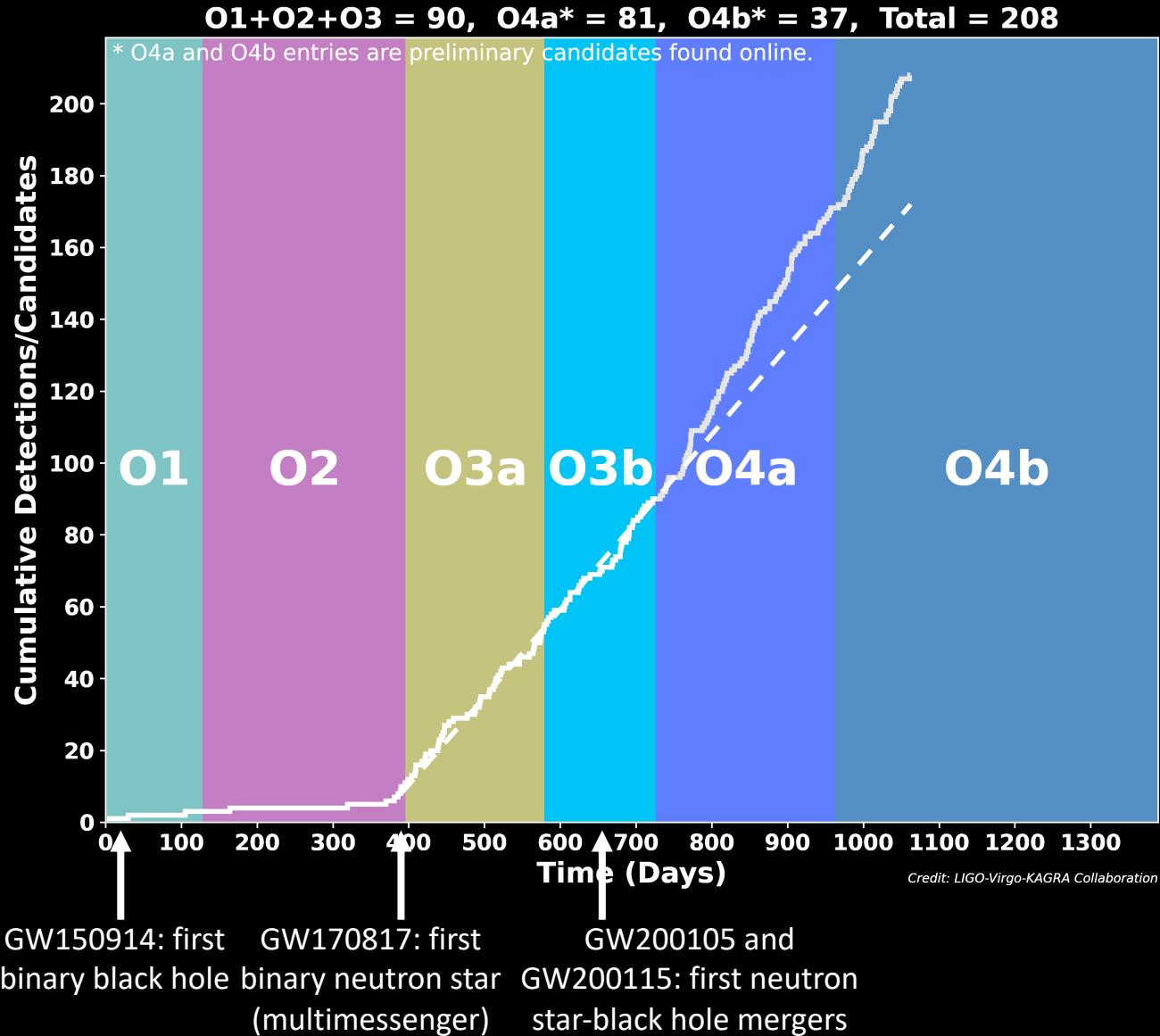
Enhances multimessenger prospects for neutron star-black hole mergers

Raises questions about X-ray binary selection effects

Backup

Our observational landscape

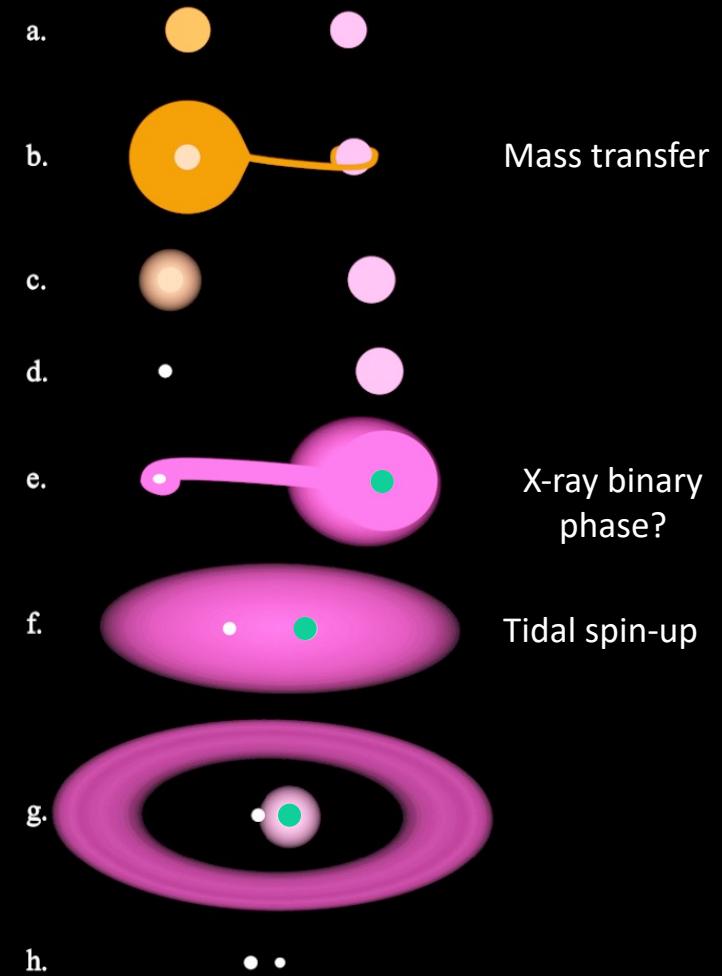




Compact-object binary formation

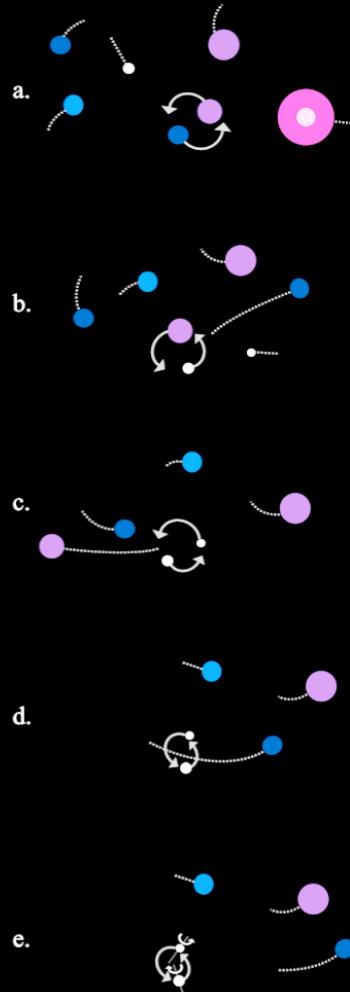
Isolated binary evolution

- Spins aligned to orbital angular momentum
- First-born black hole has small spin
- Second-born black hole can be spun up
- Masses set by supernova physics



Mandel & Farmer 1806.05820

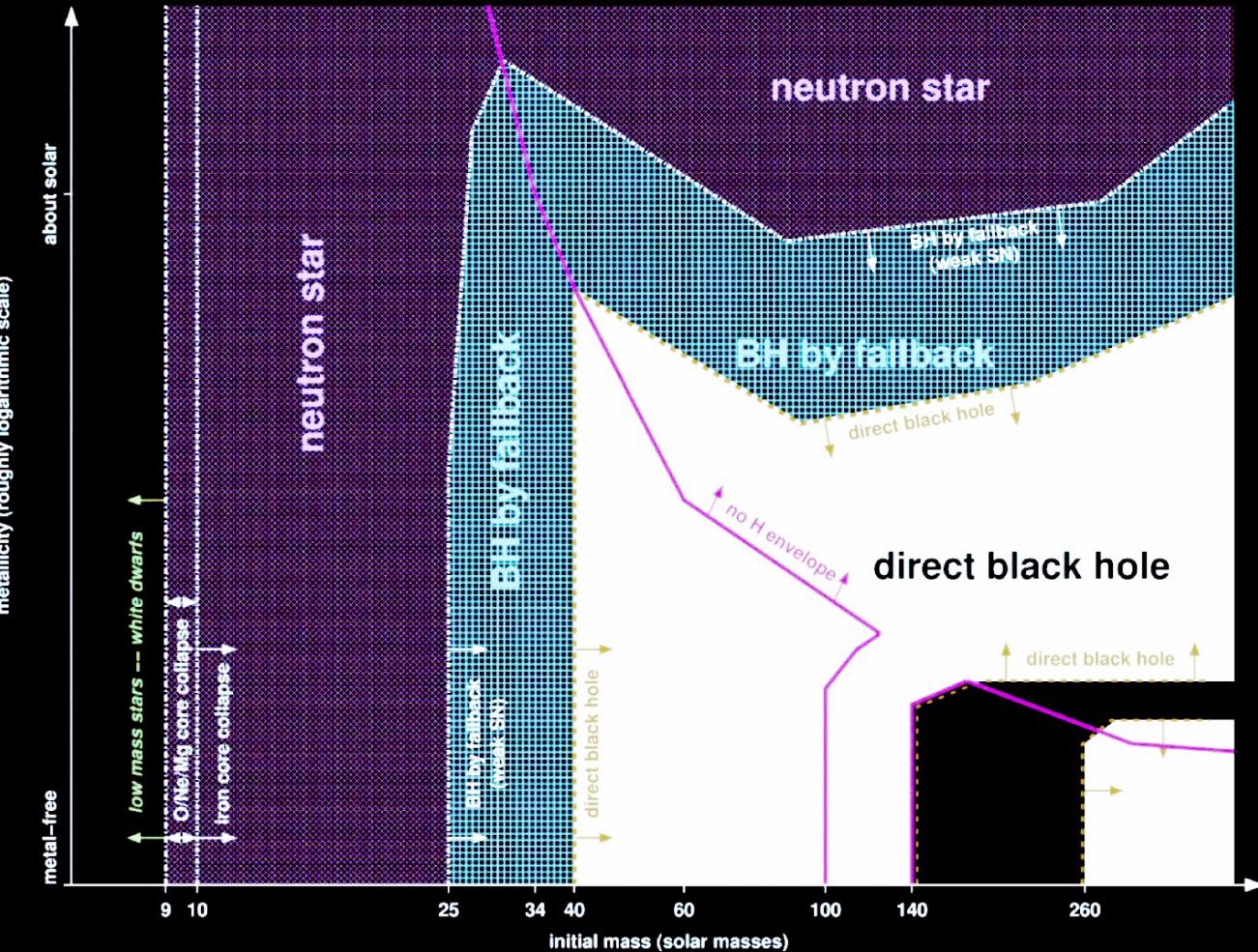
Compact-object binary formation



Dynamical assembly

- Isotropically distributed spins
- Low spins except for hierarchical mergers
- Hierarchical mergers can populate mass gaps

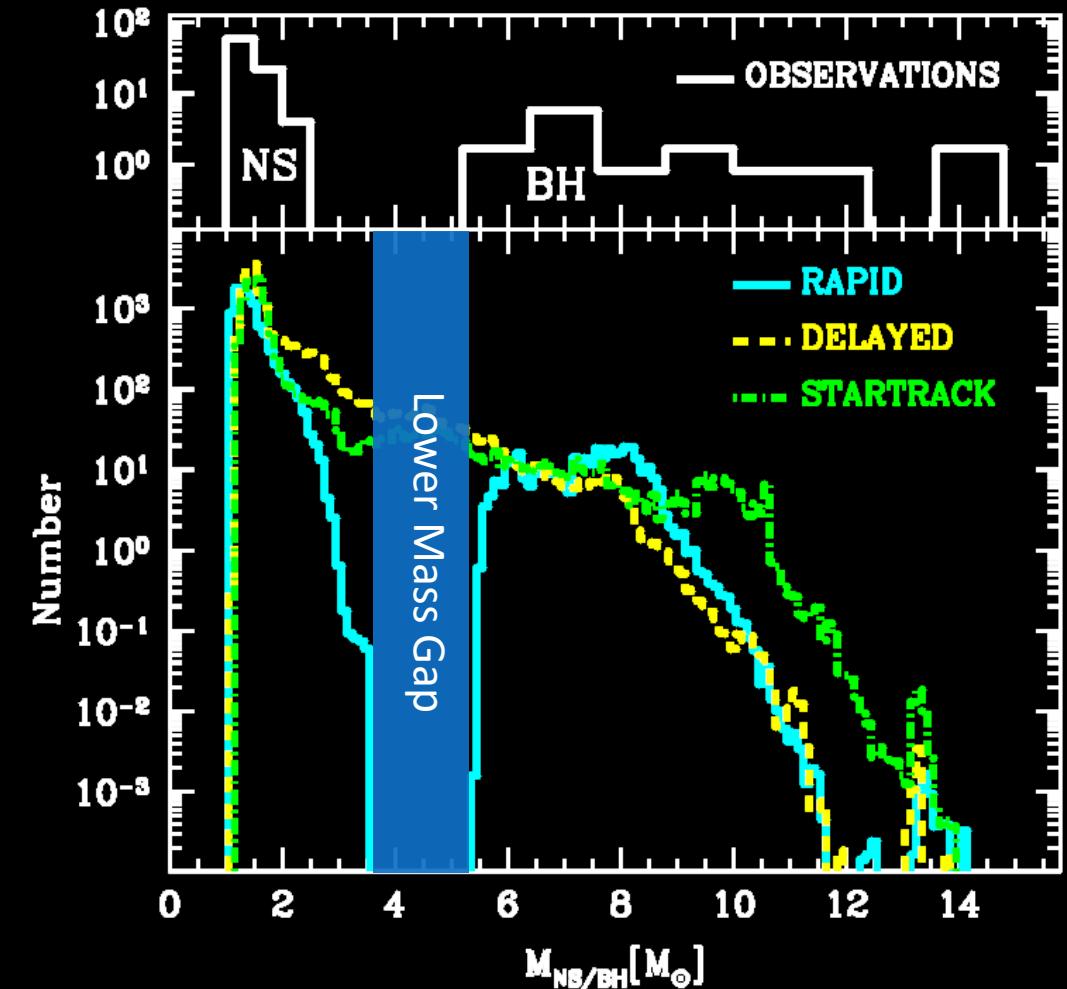
The fates of massive stars



arxiv - 2404.04248

Heger + 2003

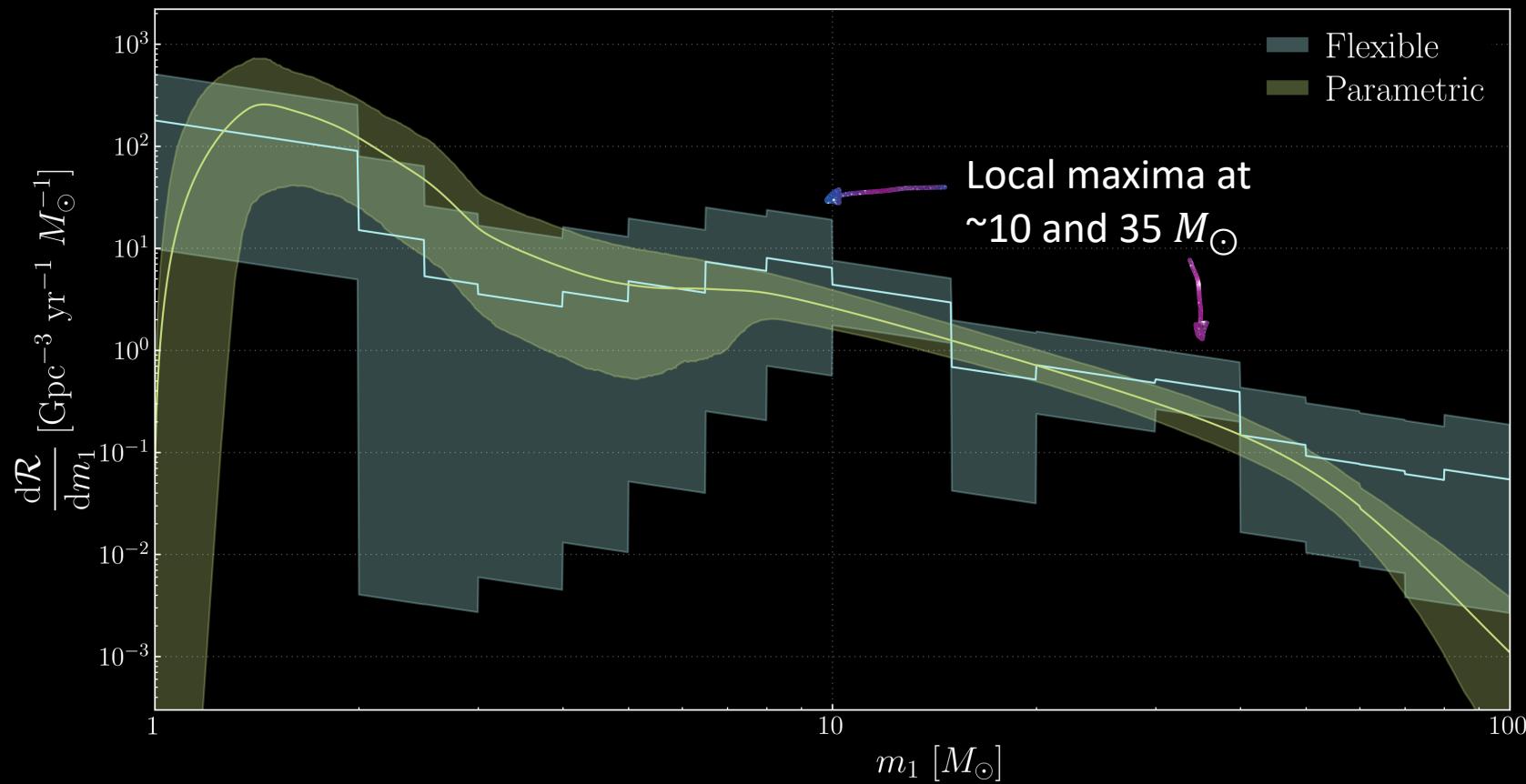
Biscoveanu - IAU 389 Plenary



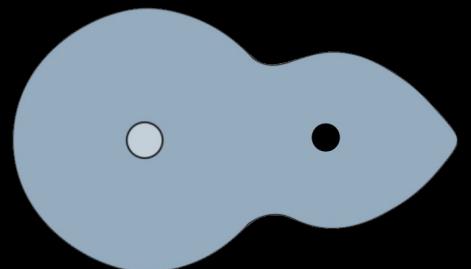
Belczynski + 2011

19

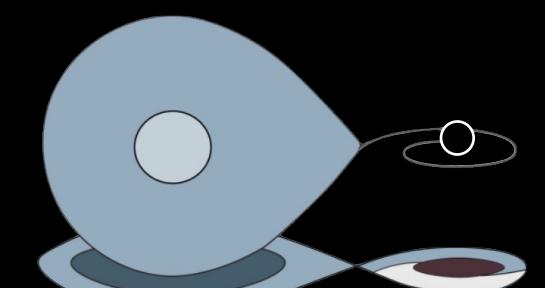
Features from different formation channels?



Common envelope



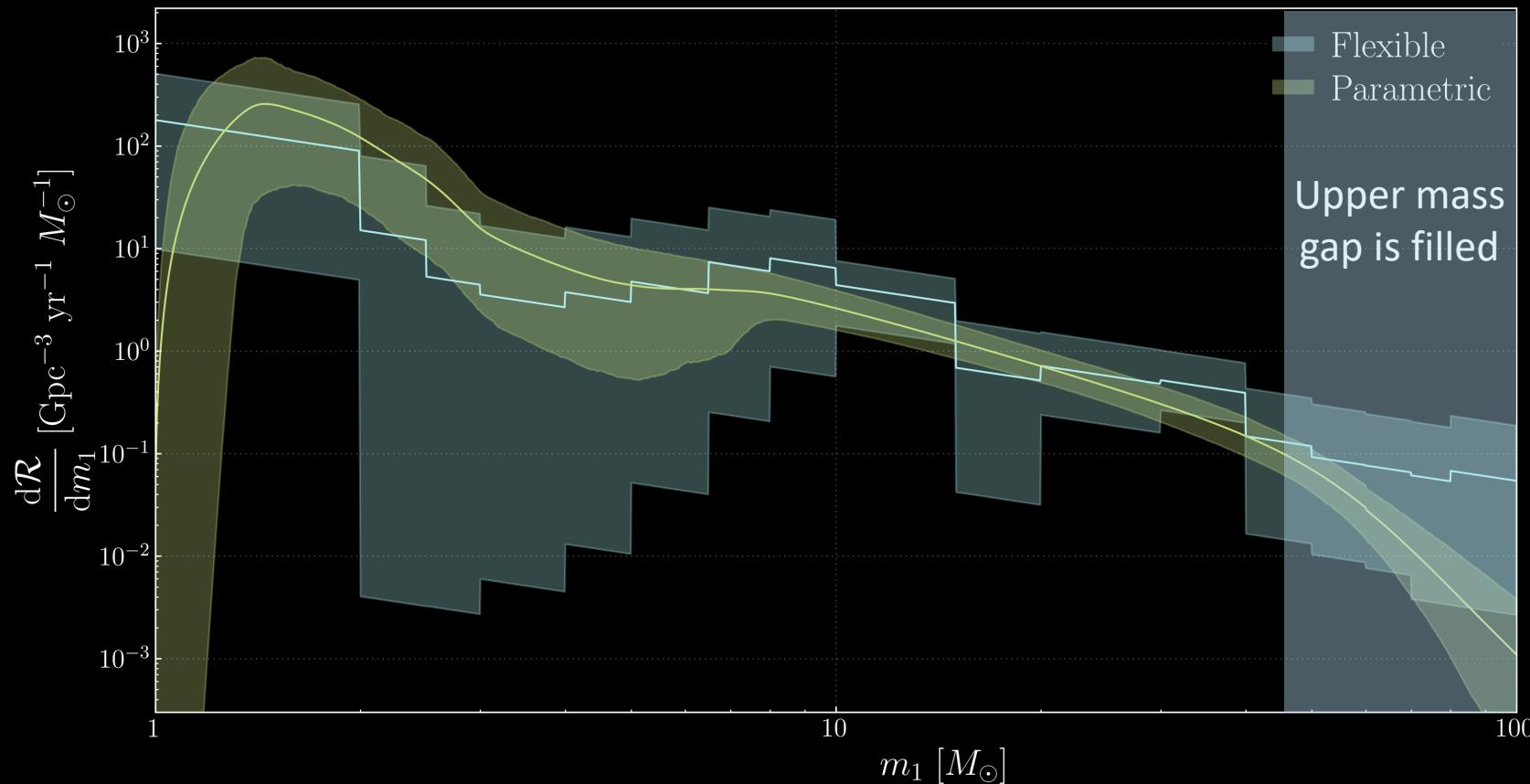
Stable mass transfer



Illustrations by Dean Kousounelos

Clues to the formation of the most massive stellar-mass black holes?

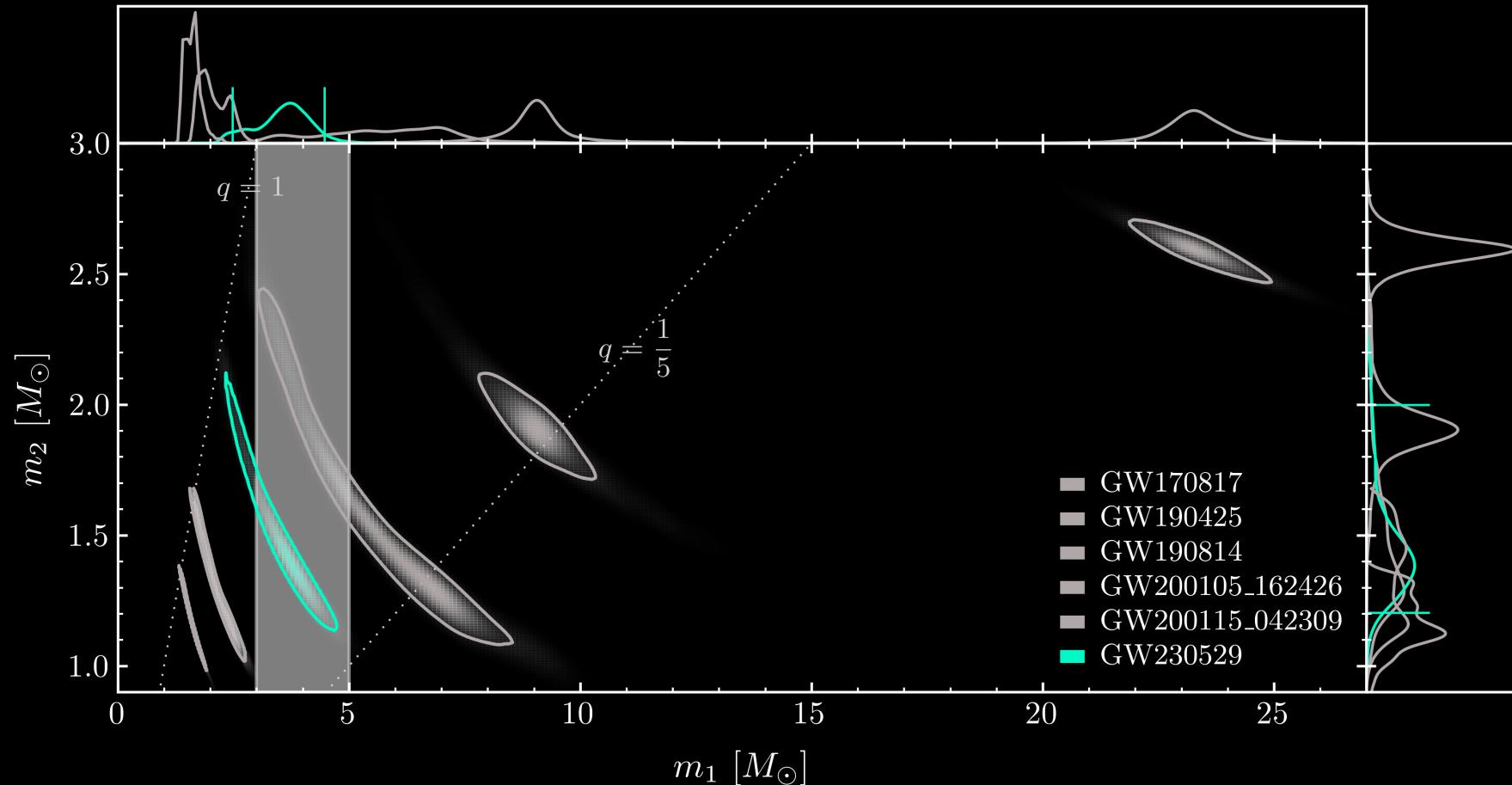
(Pulsational) pair instability supernovae



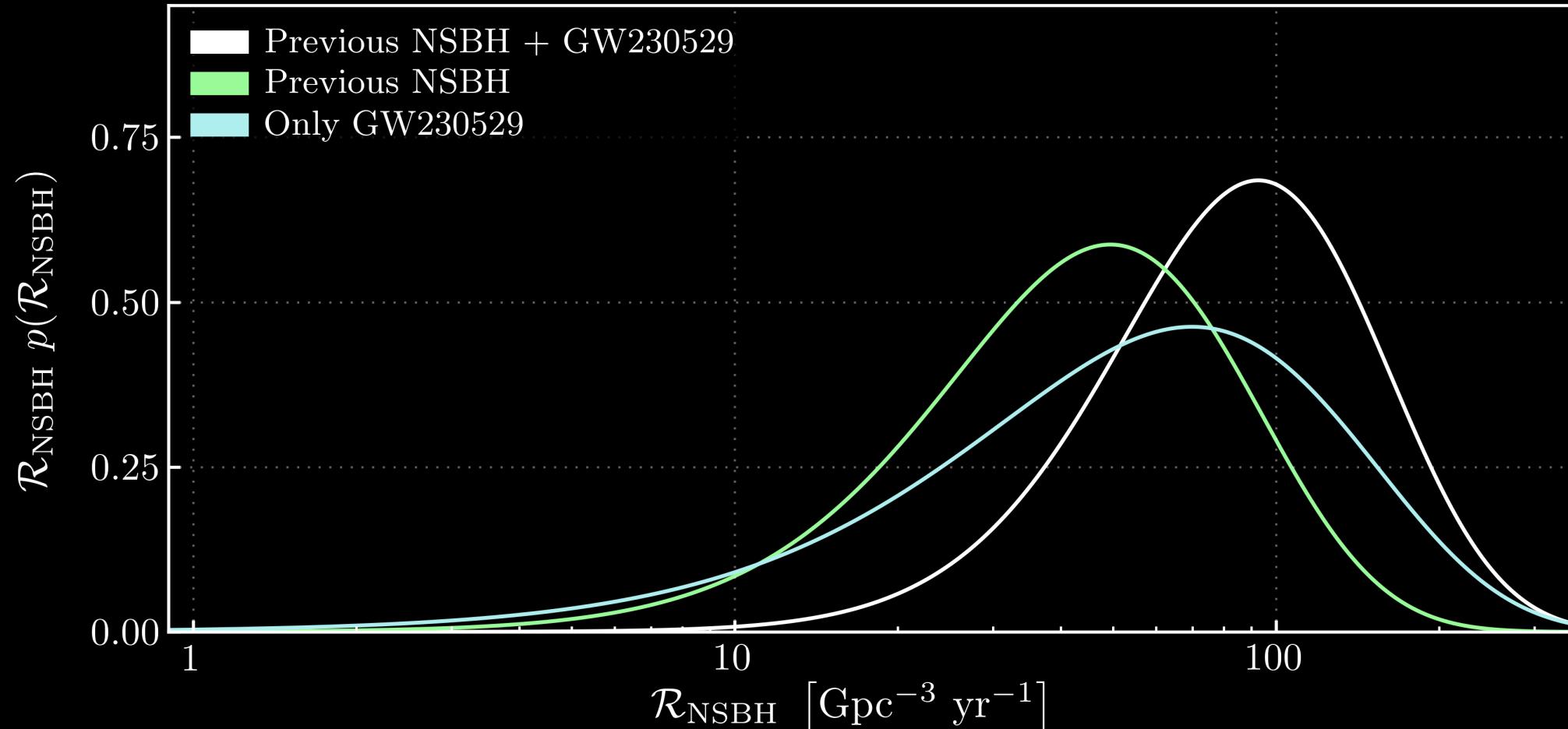
Hierarchical mergers



GW230529 is the most equal-mass neutron star-black hole merger observed to date

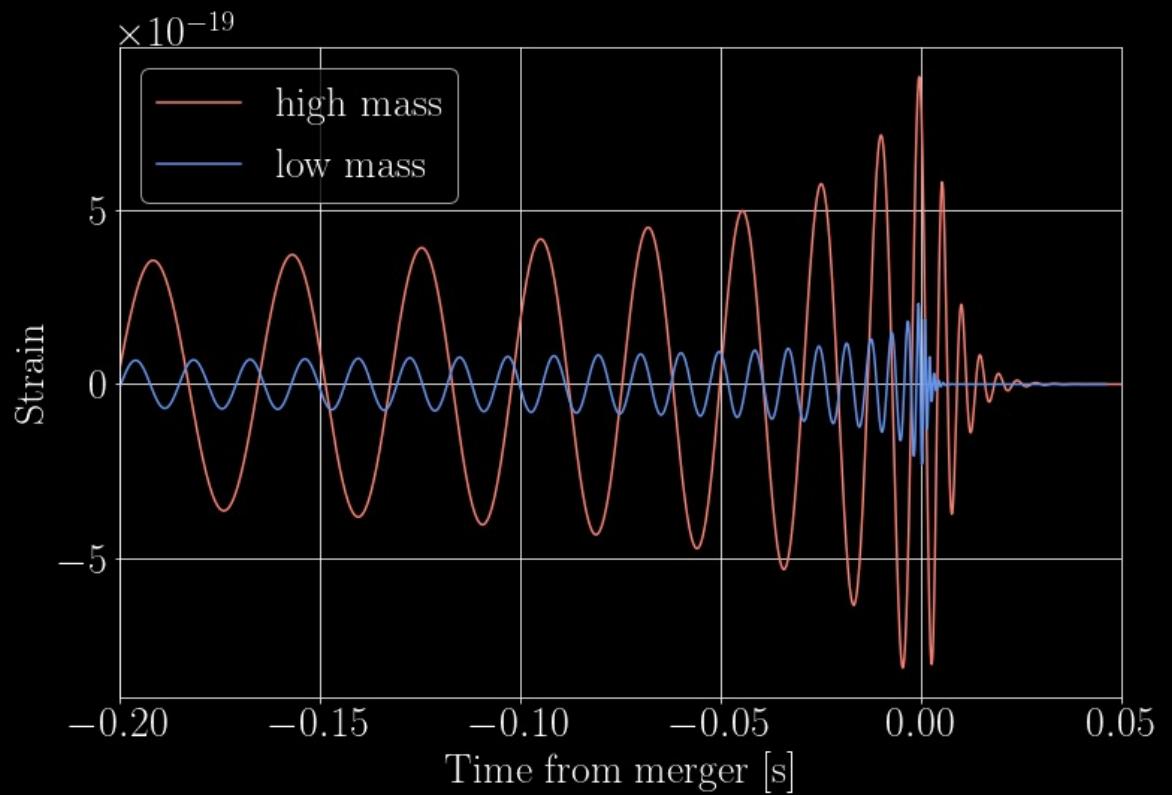


GW230529 is at least as common as more “vanilla” unequal-mass NSBH



The gravitational-wave signal

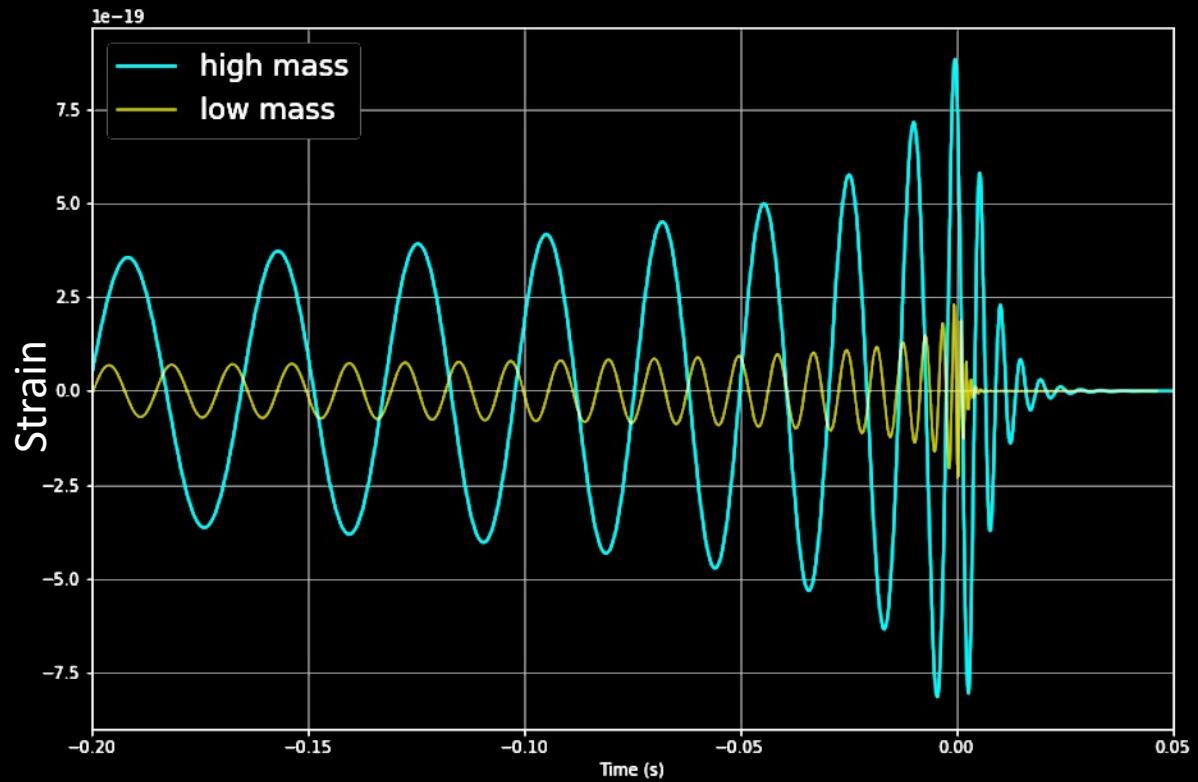
- General relativity predicts a unique gravitational waveform for each set of binary parameters
- Use different “approximant” waveforms to reduce computational cost of full solution



The gravitational-wave signal

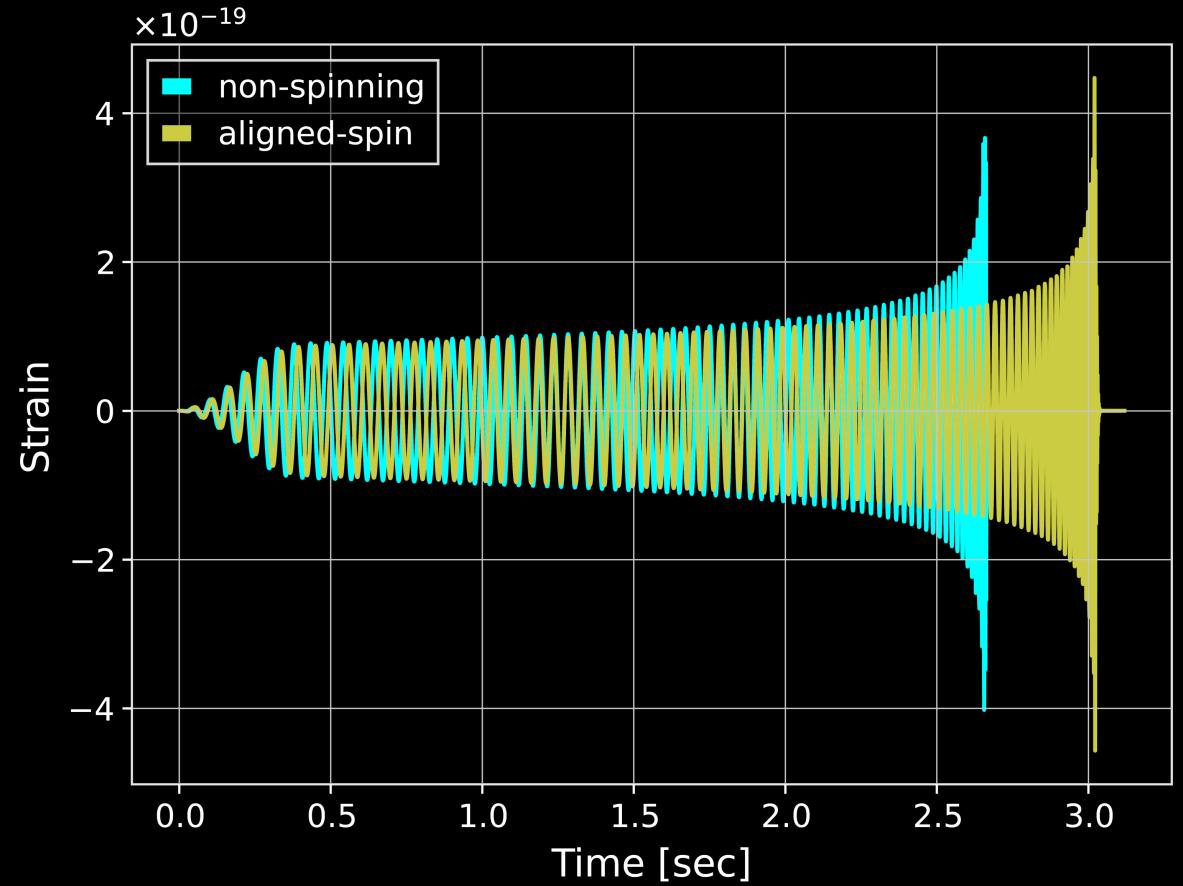
- More massive systems merge at lower frequency → spend less time in LIGO's sensitive band
- More massive, larger amplitude at fixed distance

$$\mathcal{A}_{\text{GW}} \propto \frac{\mathcal{M}^{5/6} f^{-7/6}}{d_L}$$



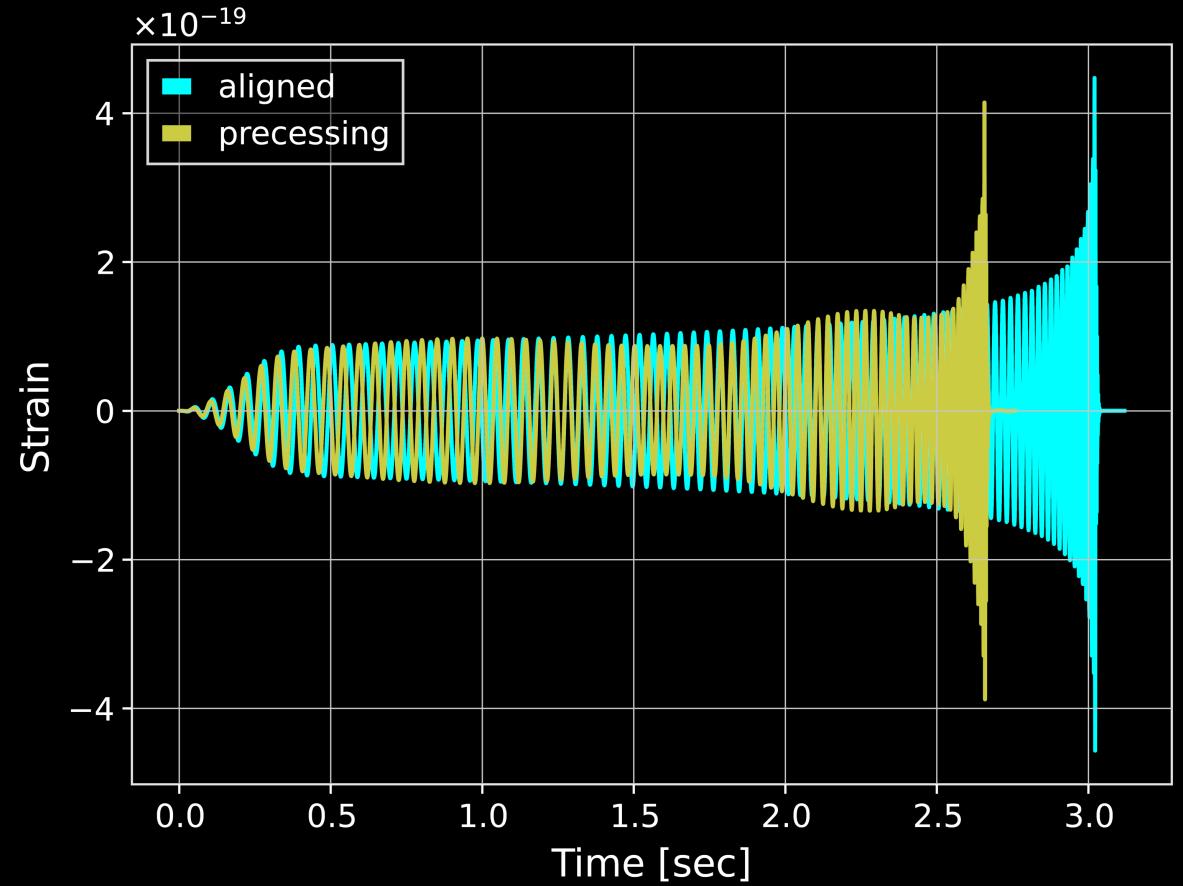
The gravitational-wave signal

- General relativity predicts a unique gravitational “waveform” for each set of binary parameters
- Amplitude and frequency both increase in time as merger approaches
- More rapidly spinning systems spend more time in LIGO’s sensitive band due to “orbital hangup effect”



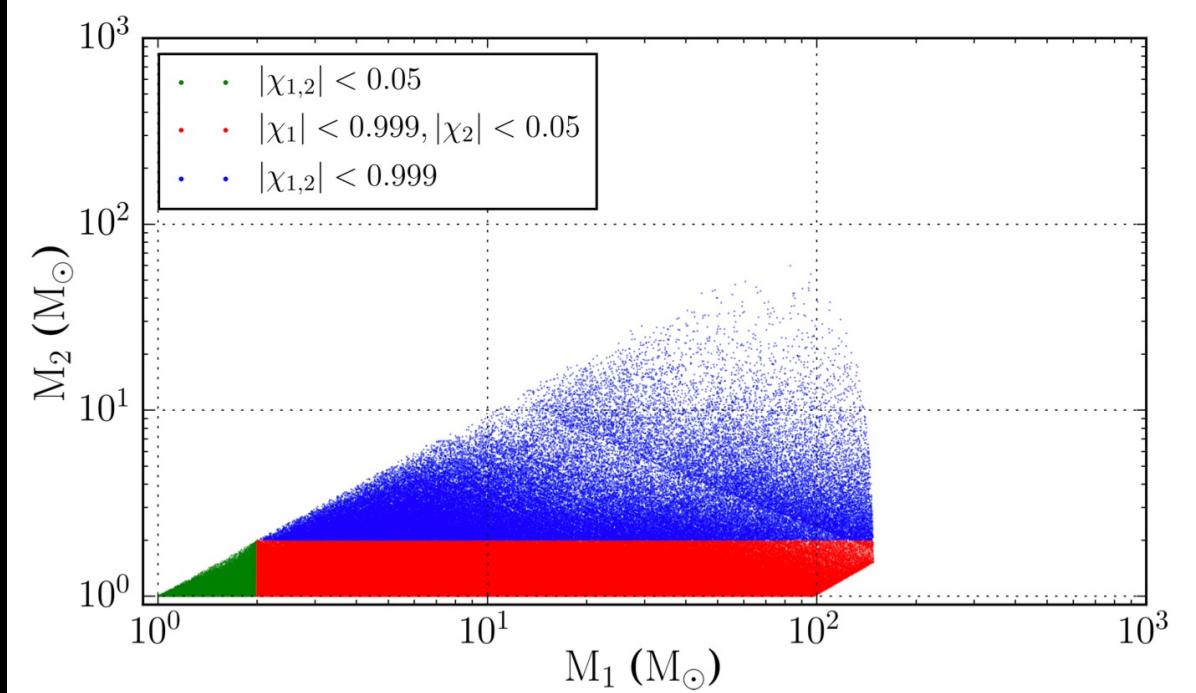
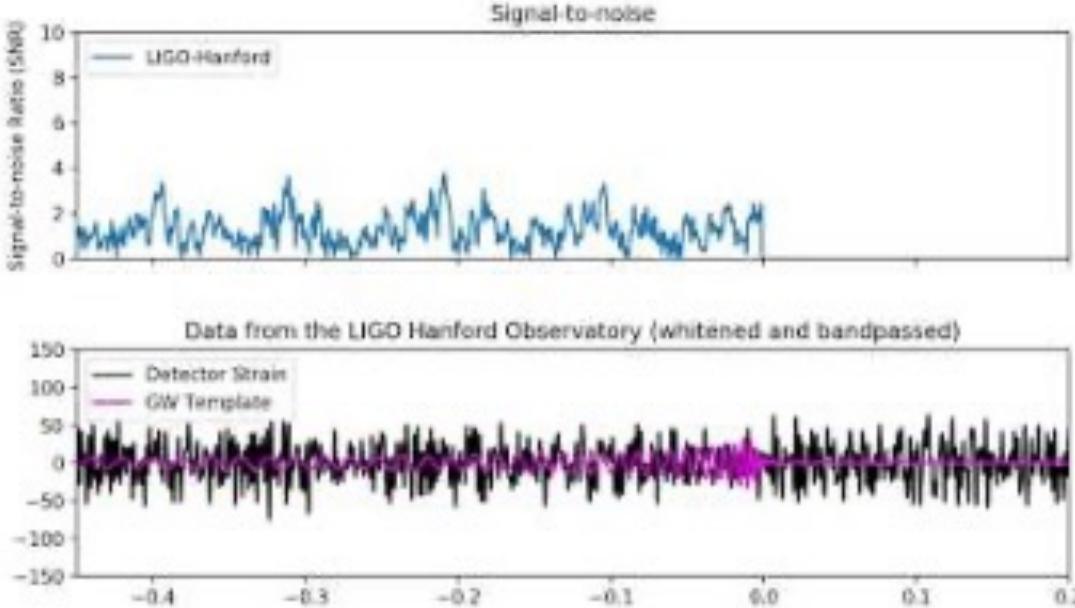
The gravitational-wave signal

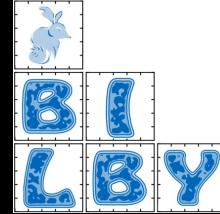
- General relativity predicts a unique gravitational “waveform” for each set of binary parameters
- Amplitude and frequency both increase in time as merger approaches
- Systems with spin vectors misaligned to the orbital angular momentum exhibit general relativistic “precession”
- Modulation of the waveform due to precession of the orbital plane



Gravitational-wave detection

- Matched filtering – compare gravitational-wave signals from a fixed template bank to the strain data, maximize the “match”





Parameter estimation

- Use Bayesian inference to sample across the whole 15-dimensional binary parameter space and obtain posterior probability distributions
- Requirements:
 - Likelihood of observing the data given a certain model
 - Prior probability of each binary parameter

$$p(\theta|d, H) = \frac{p(d|\theta, H)p(\theta|H)}{p(d|H)}$$

likelihood prior
posterior

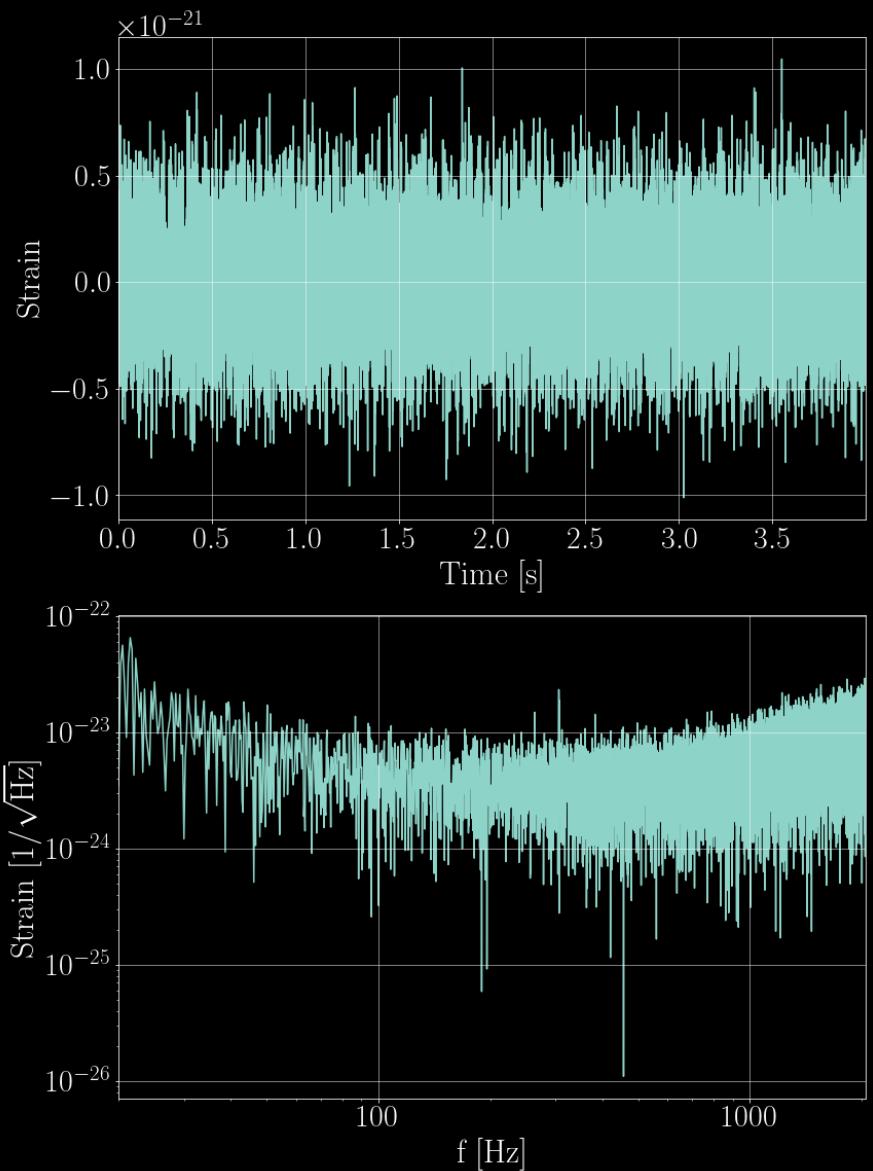
parameters data model evidence

LIGO Noise Properties

- Assume data has both a gravitational-wave signal component and a noise component:

$$d = h + n$$

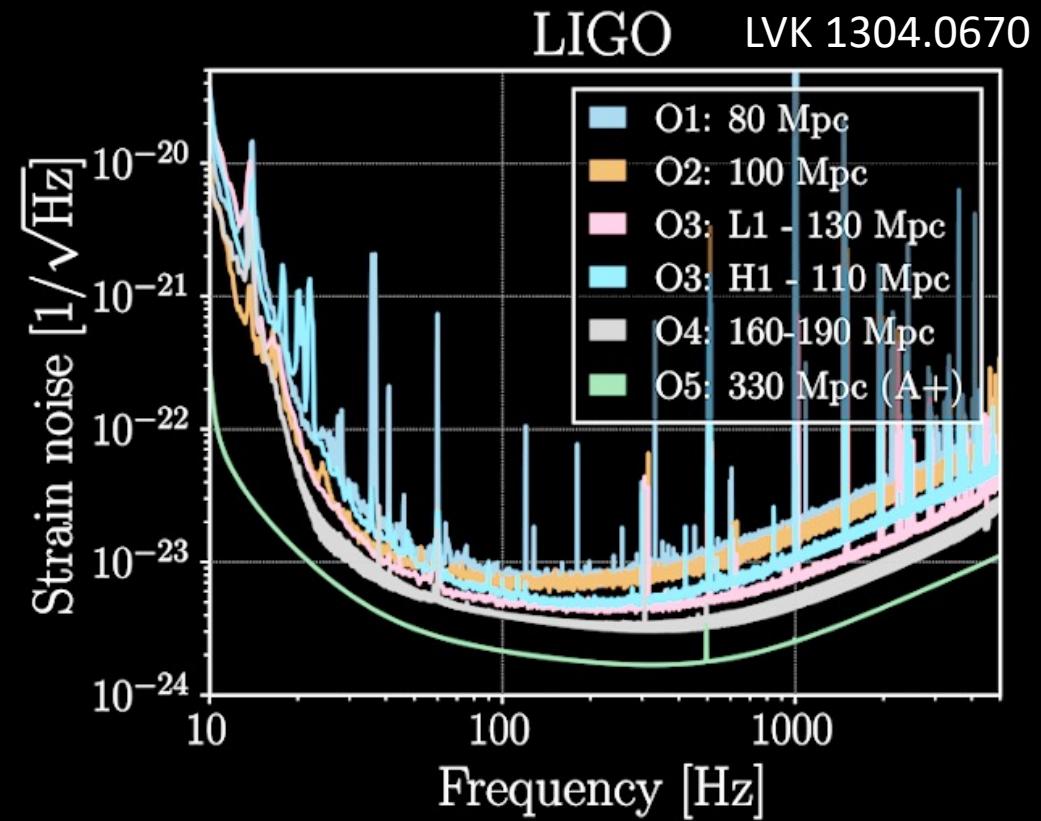
- Observe dimensionless strain time-series, Fourier transform into frequency domain



LIGO Noise Properties

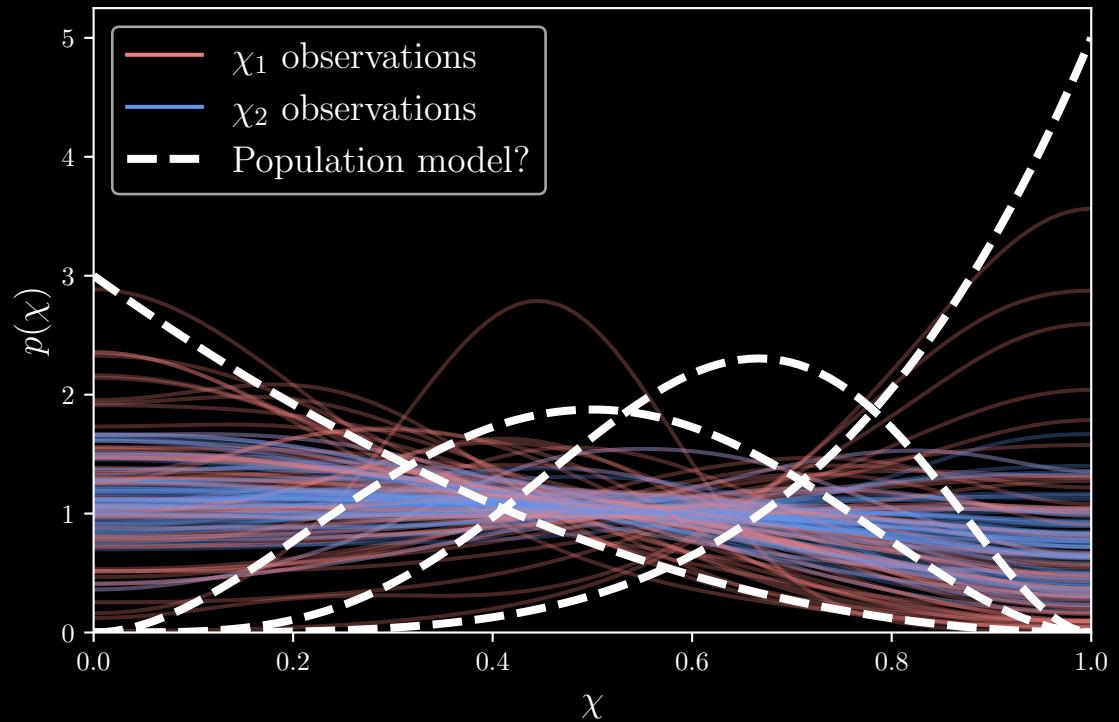
- Assume data has both a gravitational-wave signal component and a noise component:
$$d = h + n$$
- Well-behaved data in the absence of a signal is gaussian about the “amplitude spectral density” (ASD) or “strain noise” → Gaussian likelihood

$$p(d | \theta) \propto \frac{(d - h)^2}{2\text{ASD}^2}$$



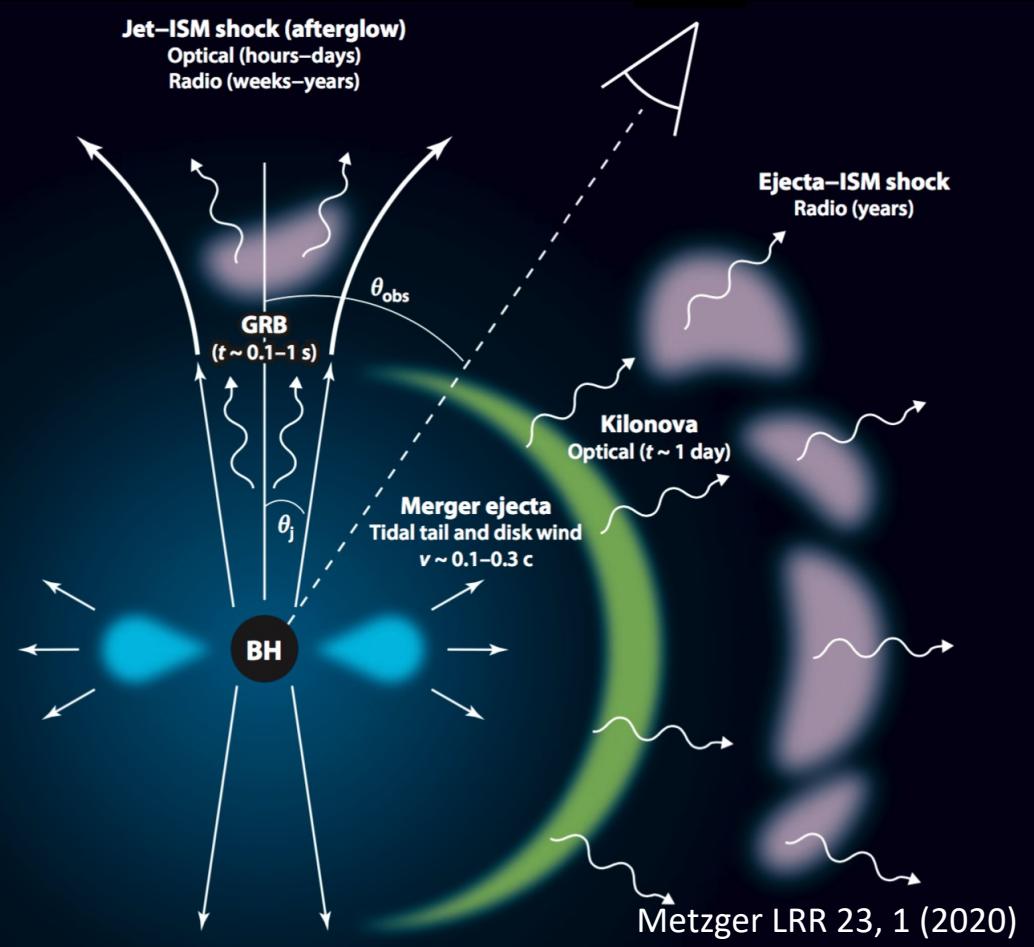
Population analysis

- By combining our $\mathcal{O}(100)$ events, we can estimate the parameters governing the distributions from which they are drawn using hierarchical Bayesian inference
- Account for individual-event statistical uncertainty and for selection effects
- Ex: assume spin magnitude is a Beta distribution characterized by an unknown mean and width
 - Wysocki+ PRD 100, 043012 (2019)



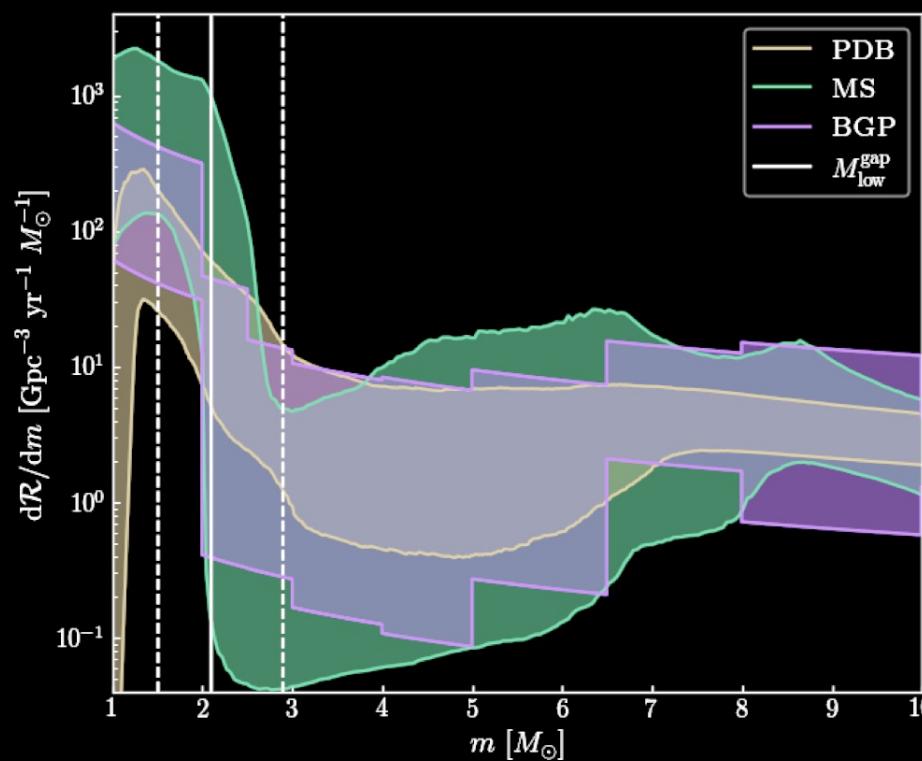
Electromagnetic counterparts

- Fate of the merger remnant and electromagnetic counterpart depends on the properties of the component objects
- *Gamma-ray burst*: most energetic electromagnetic explosions observed in the universe
- *Kilonova*: optical emission powered by the radioactive decay of freshly synthesized heavy elements

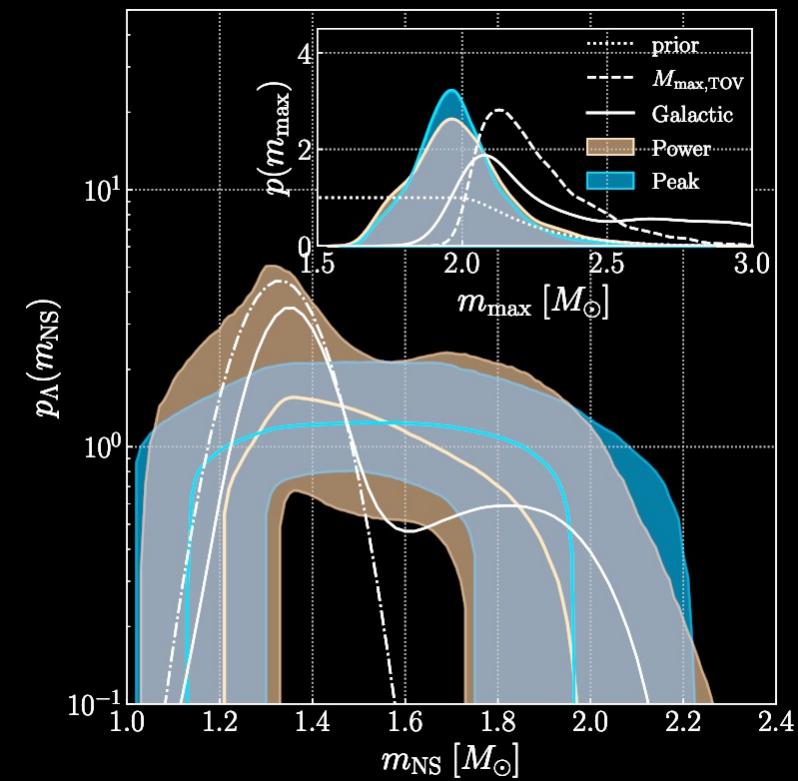


Astrophysical population inference

No evidence for lower mass gap between neutron stars and black holes



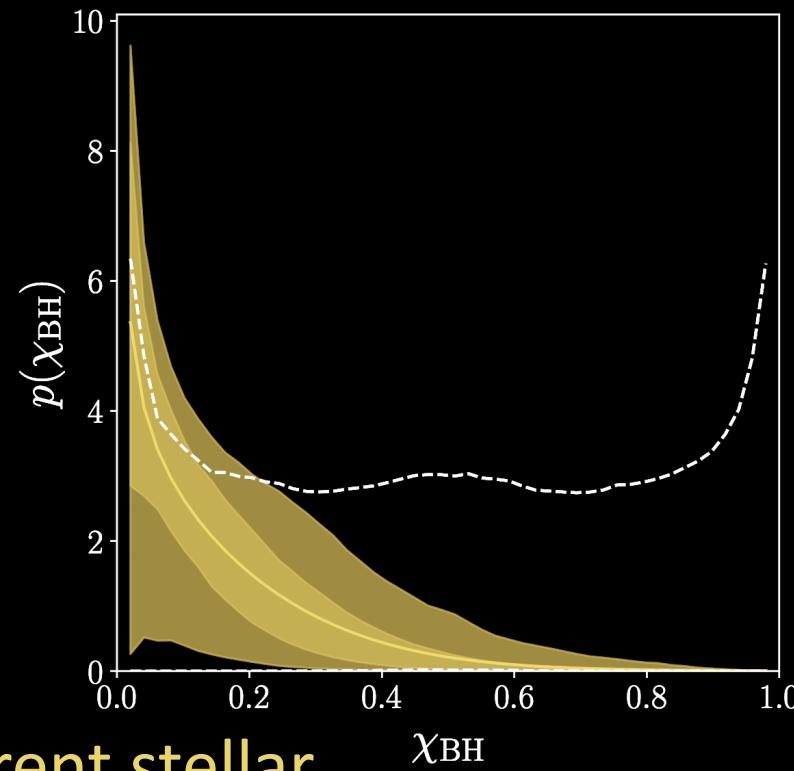
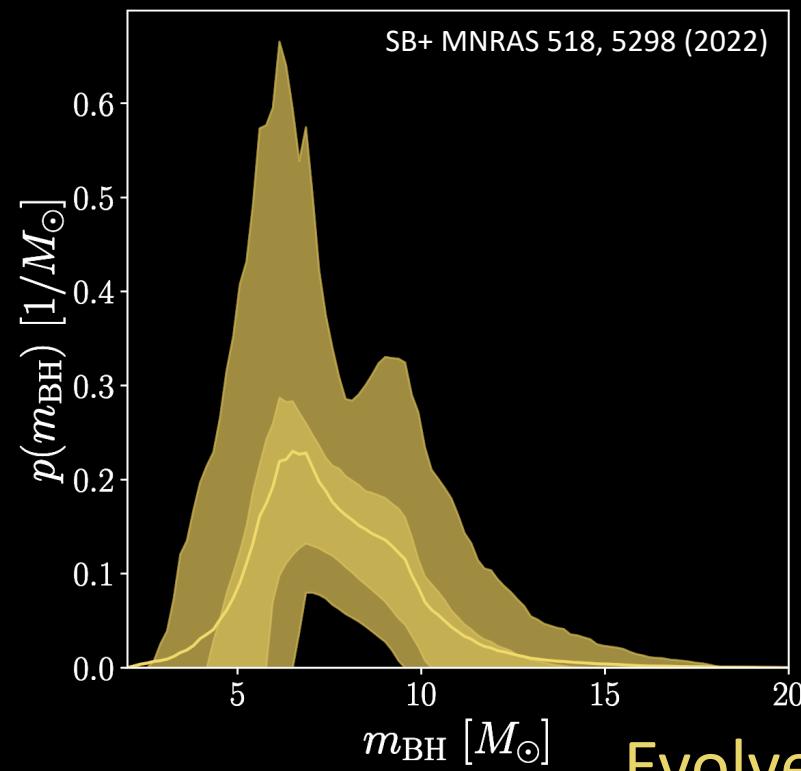
Neutron stars detected in GWs are more massive than those detected as pulsars in our galaxy



Neutron star-black hole mergers

The black holes are less massive than those in black hole binaries

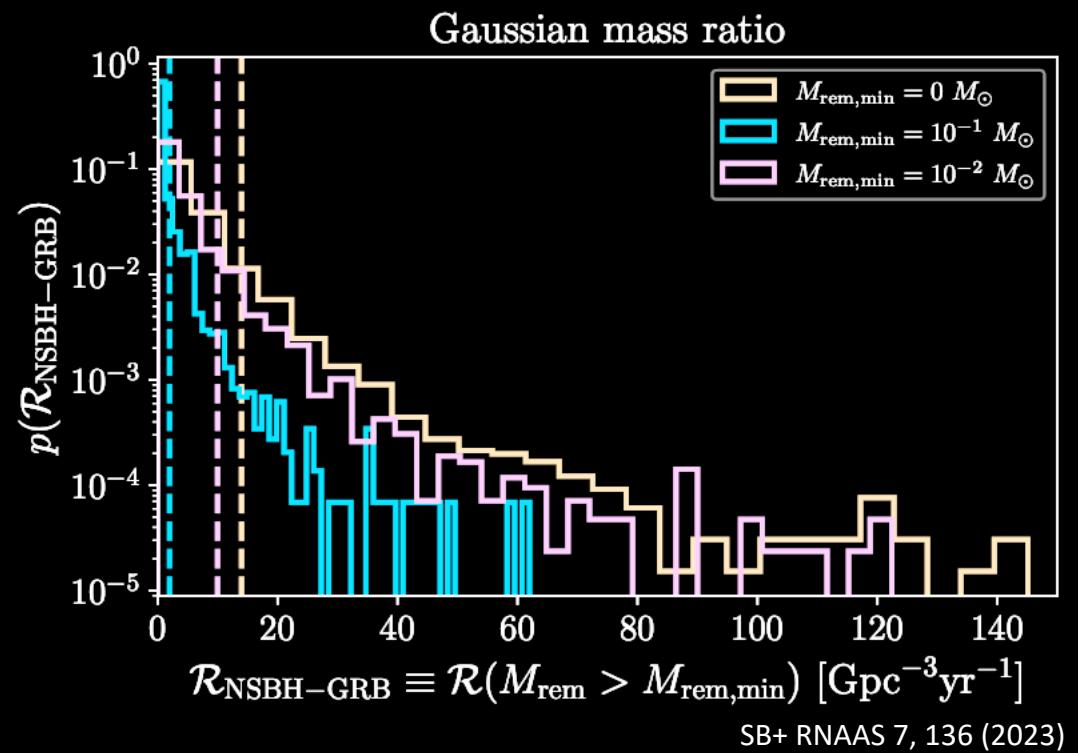
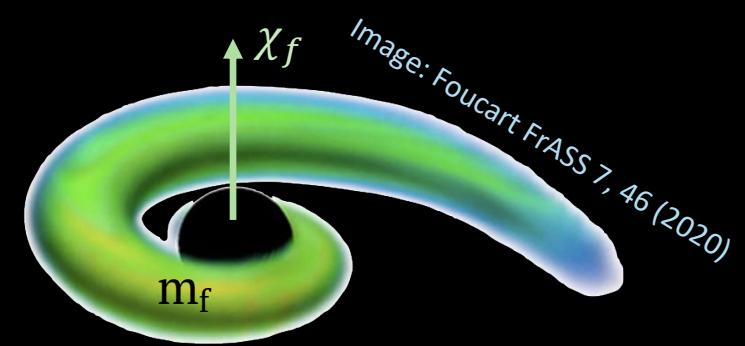
And more slowly spinning



Evolve from different stellar progenitors

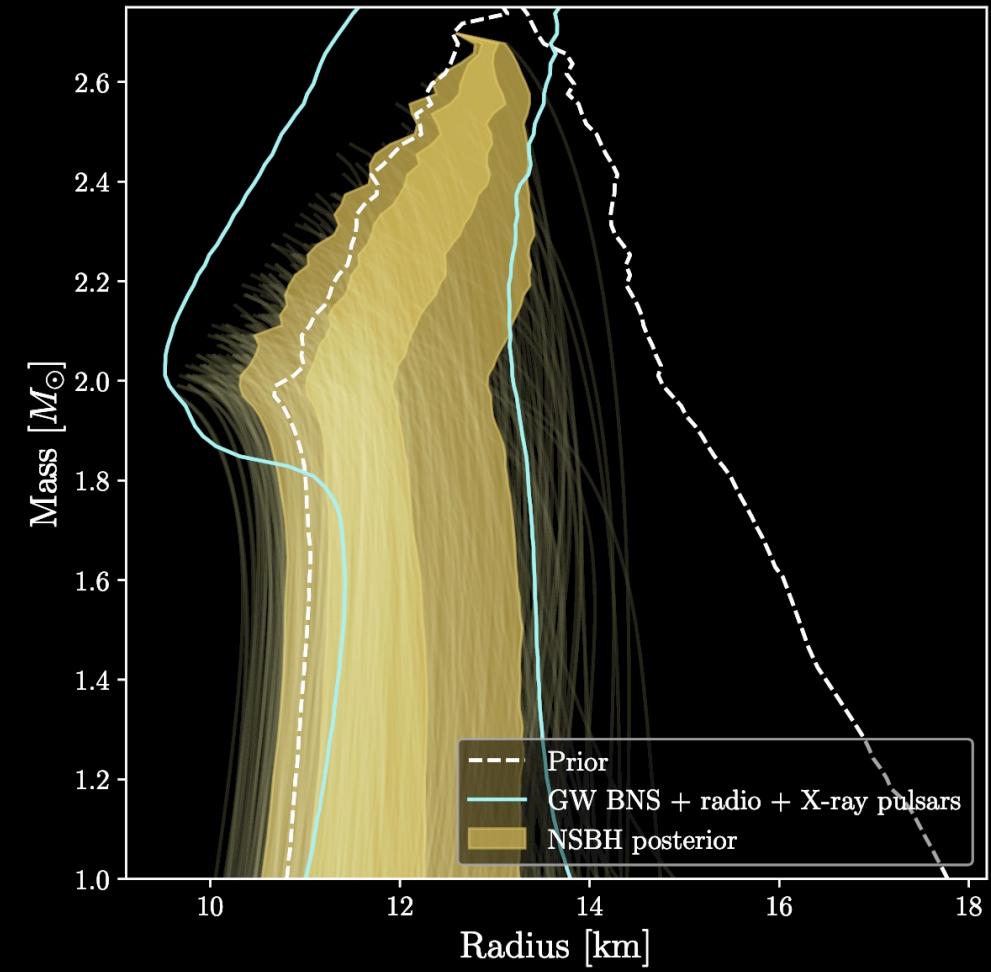
Multimessenger prospects

- Neutron-star black hole mergers are unlikely to produce electromagnetic counterparts
 - Unequal mass ratios and low black hole spins
- Upper limit of $20 \text{ Gpc}^{-3} \text{ yr}^{-1}$ GRBs from NSBH mergers at 90% credibility (underlying beaming-corrected rate)



Neutron star-black hole mergers

- For small neutron star radii, extreme mass ratios, and low black hole spin, the neutron star is more likely to plunge directly into the black hole → no electromagnetic counterpart
- Based on the observed population, at most 14% of mergers detectable in gravitational waves may be electromagnetically bright
- Fold in counterpart non-detection into Bayesian framework to constrain the neutron star equation of state



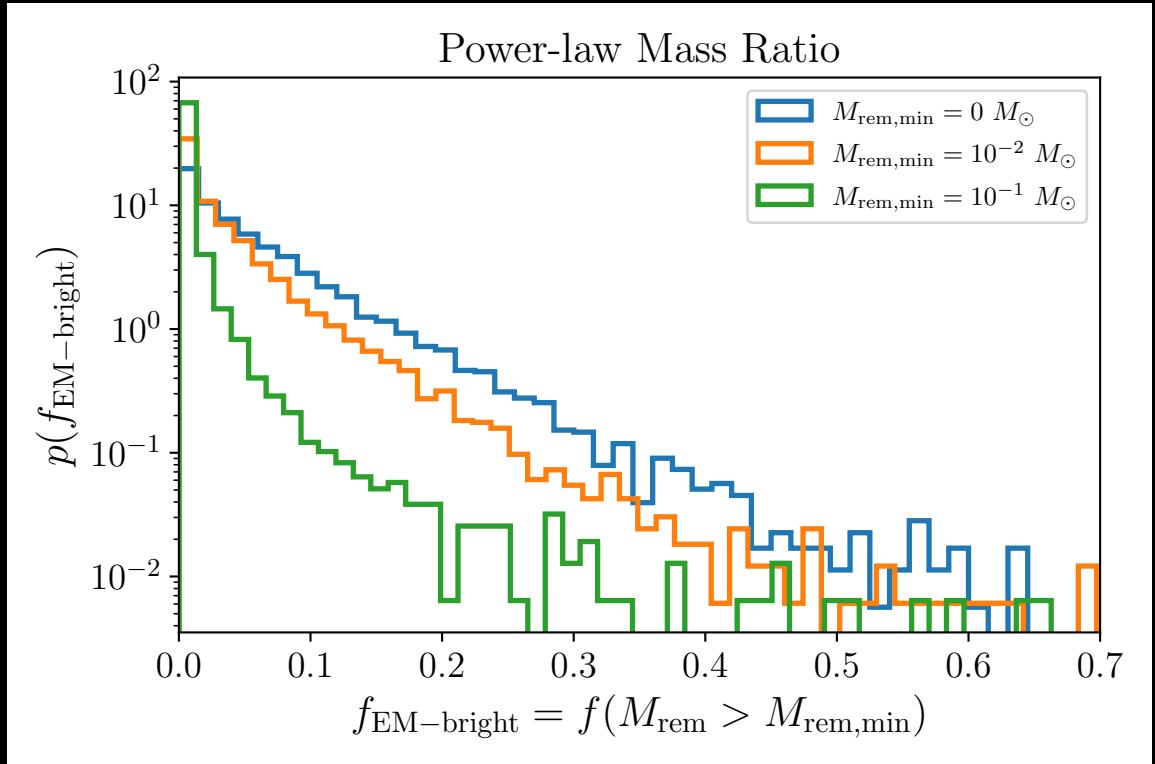
Electromagnetically-bright fraction

- At most 14% of NSBH sources may have an associated electromagnetic counterpart
- 99% probability that none of the four NSBH candidates we consider were electromagnetically bright
 - Use this constraint to inform the population properties and measure the neutron star equation of state

$$\text{observed remnant mass} = 0$$
$$p(M_{\text{rem}} | q, m_{\text{NS}}, \chi_{\text{BH},z}, \Lambda_{\text{EoS}}) = \delta(\hat{M}_{\text{rem}}(q, m_{\text{NS}}, \chi_{\text{NS}}, \chi_{\text{BH},z}, \Lambda_{\text{EoS}}))$$

remnant mass predicted by fitting formula

equation of state
piecewise polytrope
parameters



Hierarchical Modeling

- The new likelihood is the original likelihood **marginalized** over the original parameters:

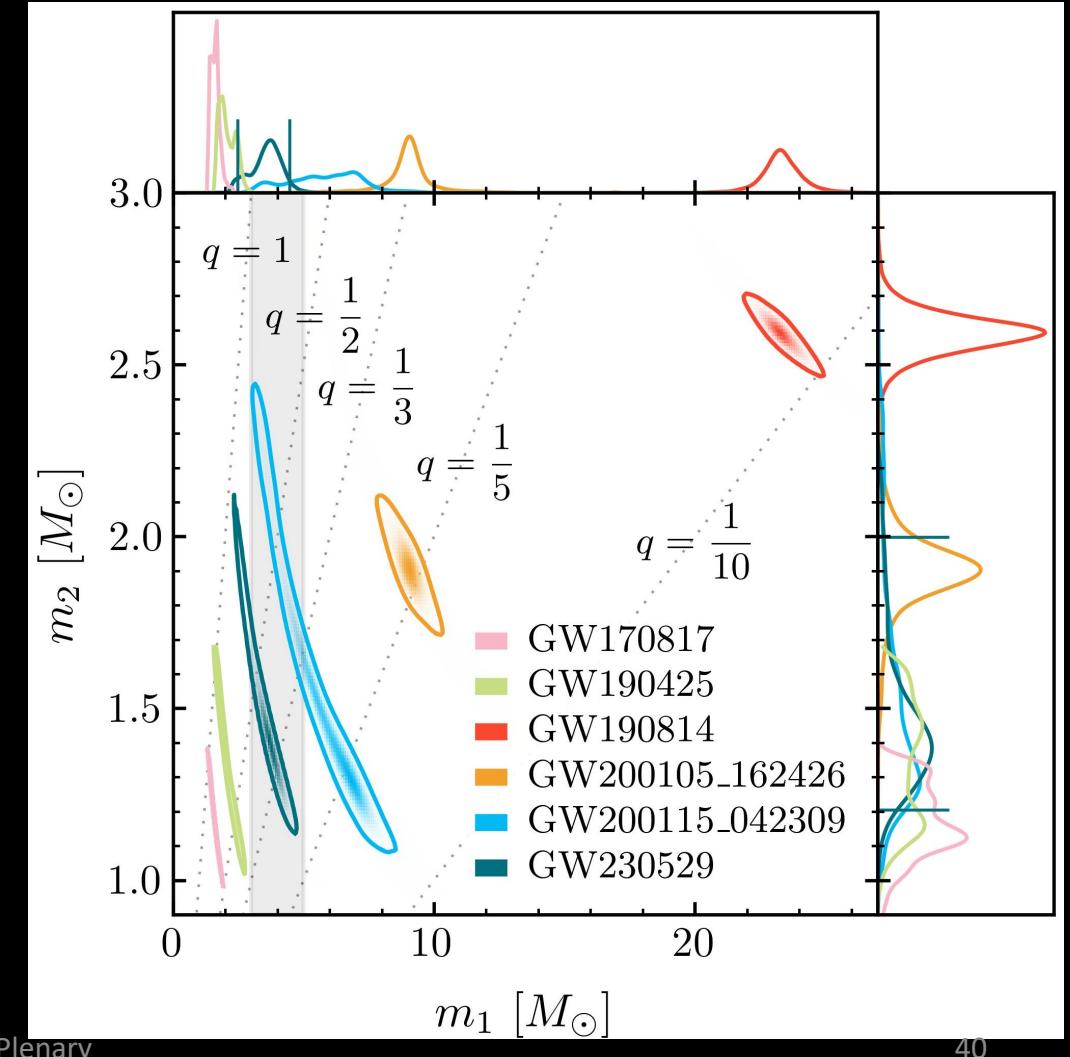
$$\begin{aligned}\mathcal{L}(d|\Lambda) &= \int d\theta \mathcal{L}(d|\theta, \Lambda) \pi(\theta|\Lambda) && \text{Hyper-prior} \\ &= \int d\theta \mathcal{L}(d|\theta) \pi(\theta|\Lambda) && \text{Original likelihood (doesn't depend on hyper-parameters)} \\ &= \int d\theta \frac{p(\theta|d) Z_\theta}{\pi_0(\theta)} \pi(\theta|\Lambda) && \text{Original evidence} \\ &&& \text{Original prior}\end{aligned}$$

GW230529

- Observed by LIGO Livingston only – significant detection by three independent search pipelines

	GstLAL	MBTA	PyCBC
Online S/N	11.3	11.4	11.6
Online inverse FAR (yr)	1.1	1.1	160.4
Offline inverse FAR (yr)	60.3	>1000	>1000

- Mass of the primary staunchly in $3\text{-}5 M_{\odot}$ range
- Support for either very high-mass neutron star or low-mass black hole



What's next?

