

# The Origin of Universality in the Inner Edges of Planetary Systems



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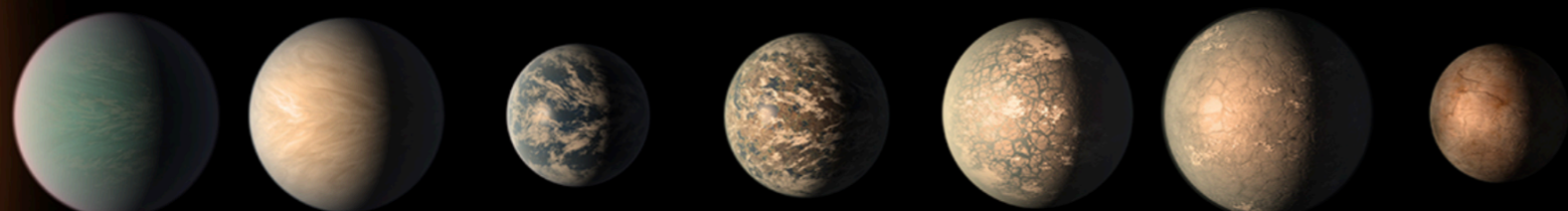
**Juliette Becker (U Wisconsin-Madison)**

Jupiter ( $M \approx 0.001 M_{\text{sun}}$ )



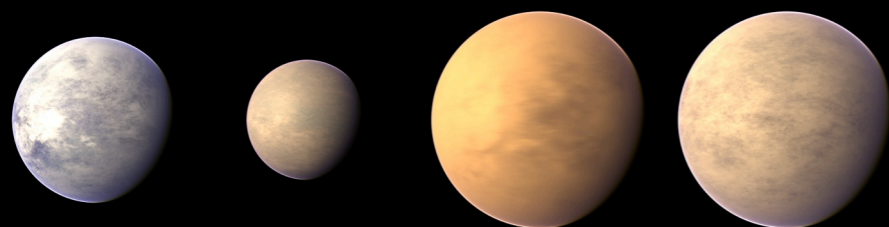
$P_{\text{Io}} \approx 1.8$  days

Trappist 1 ( $M \approx 0.09 M_{\text{sun}}$ )

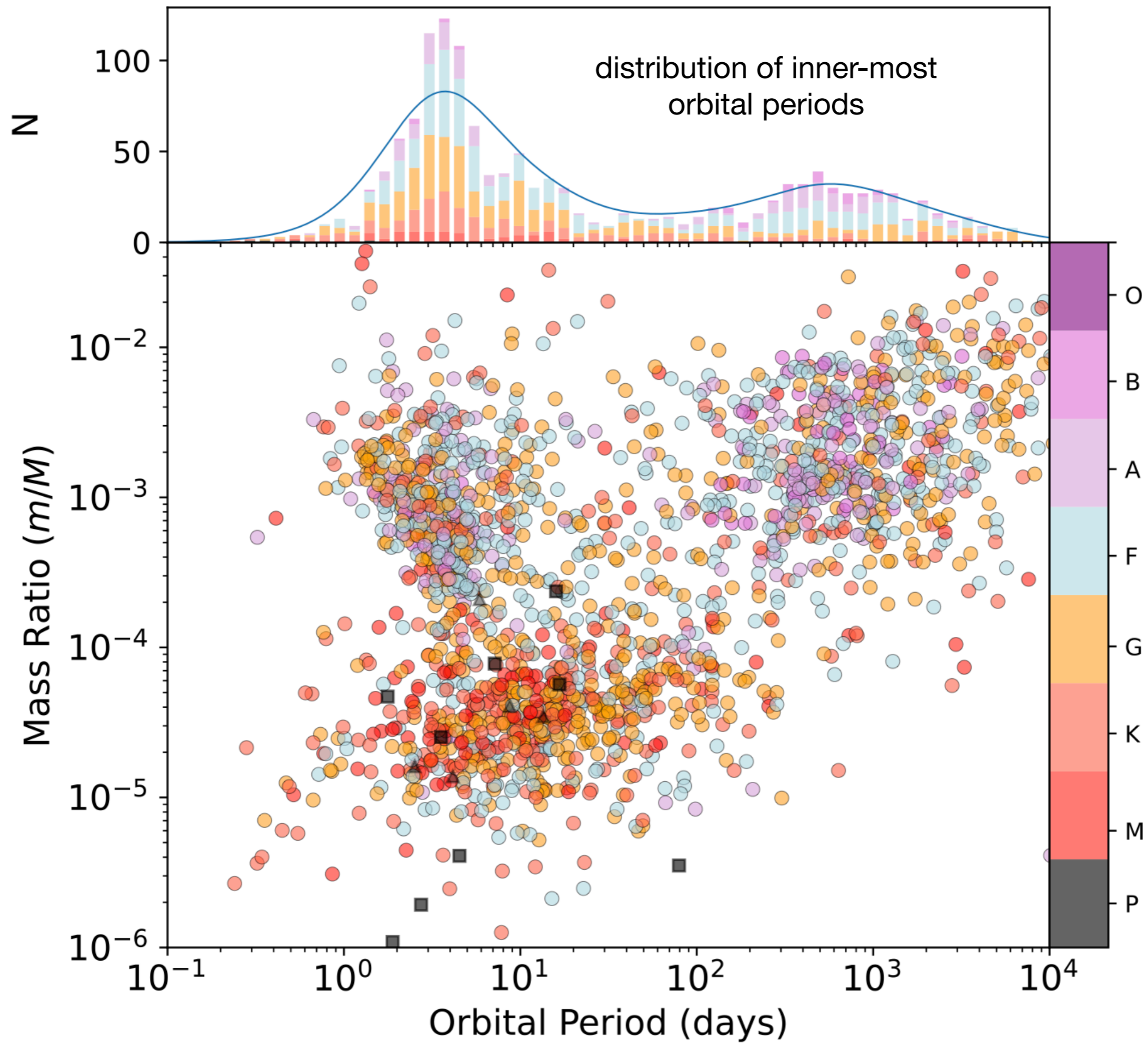


$P_{\text{b}} \approx 1.5$  days

Kepler 256 ( $M \approx 1 M_{\text{sun}}$ )



$P_{\text{b}} \approx 1.6$  days

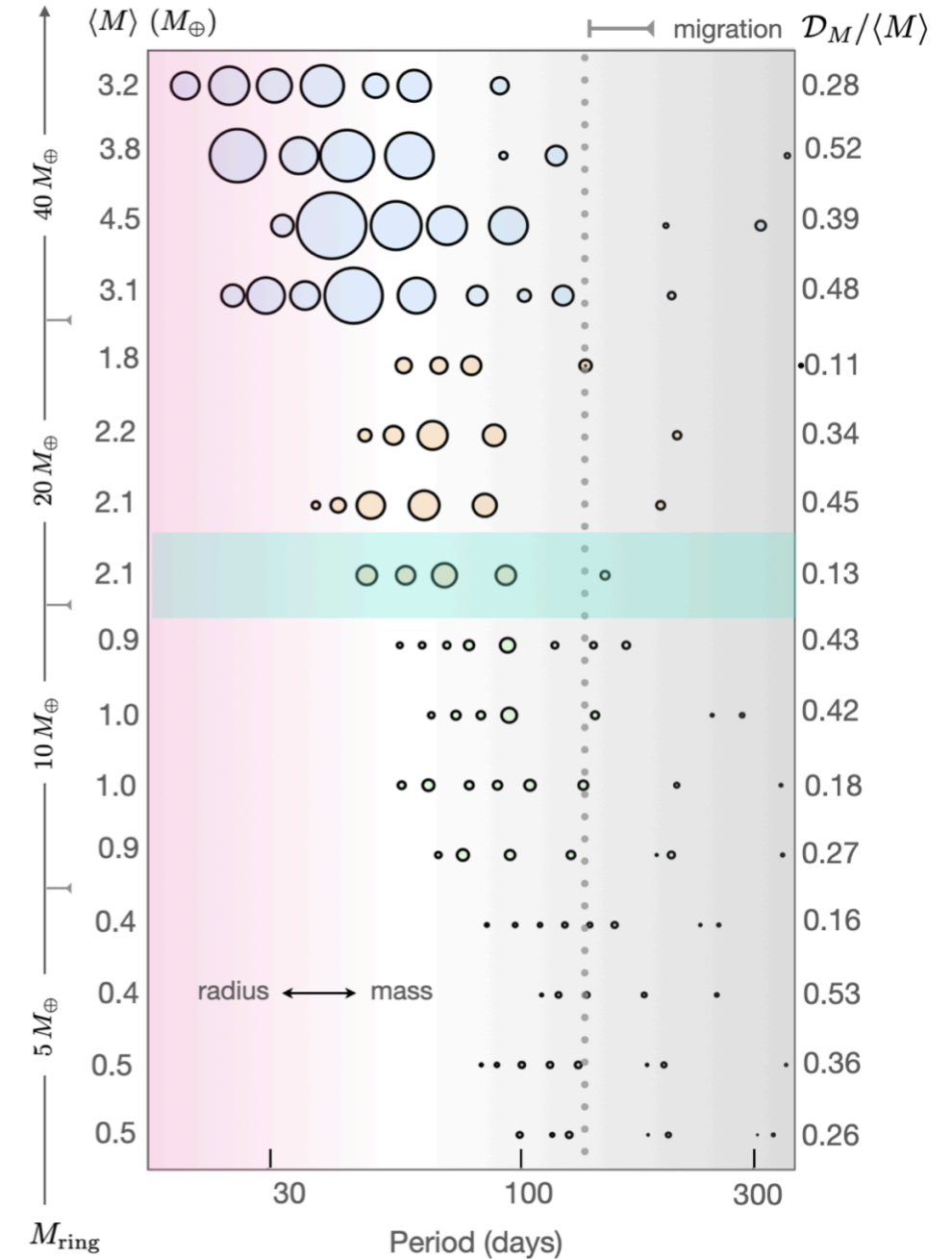
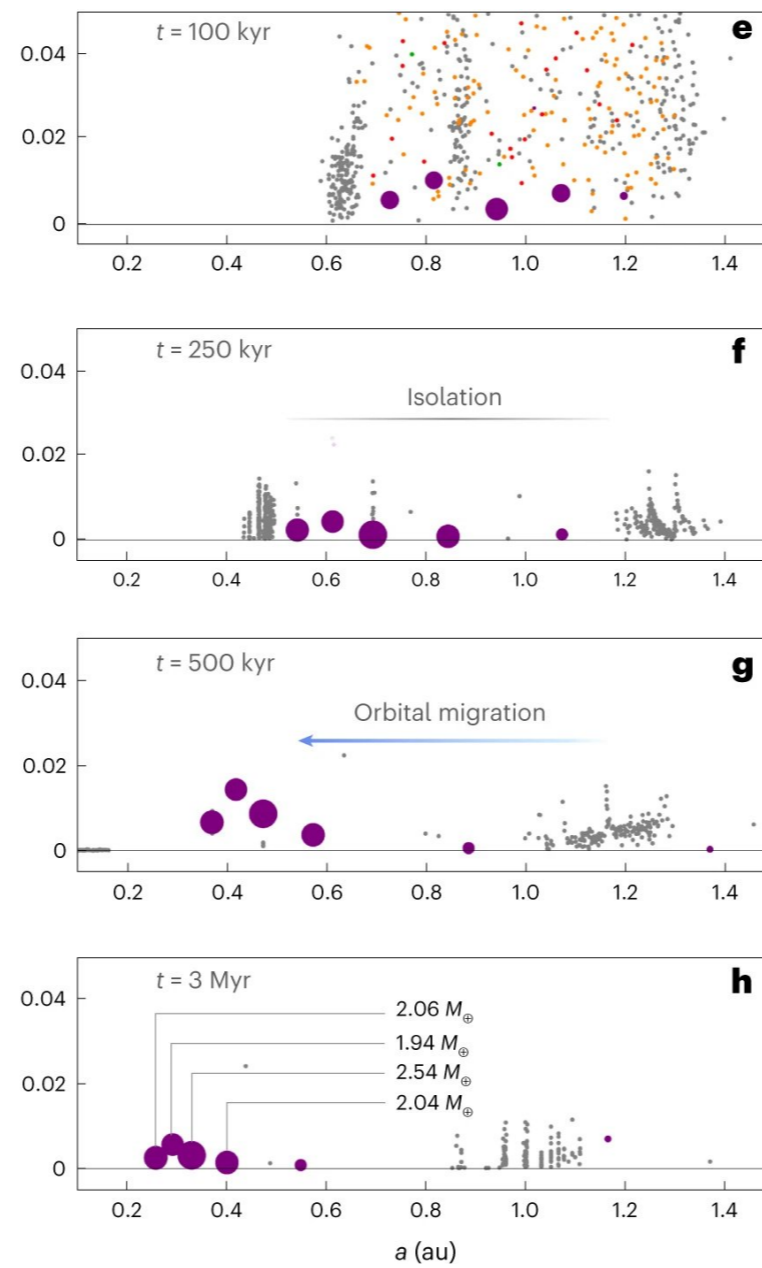
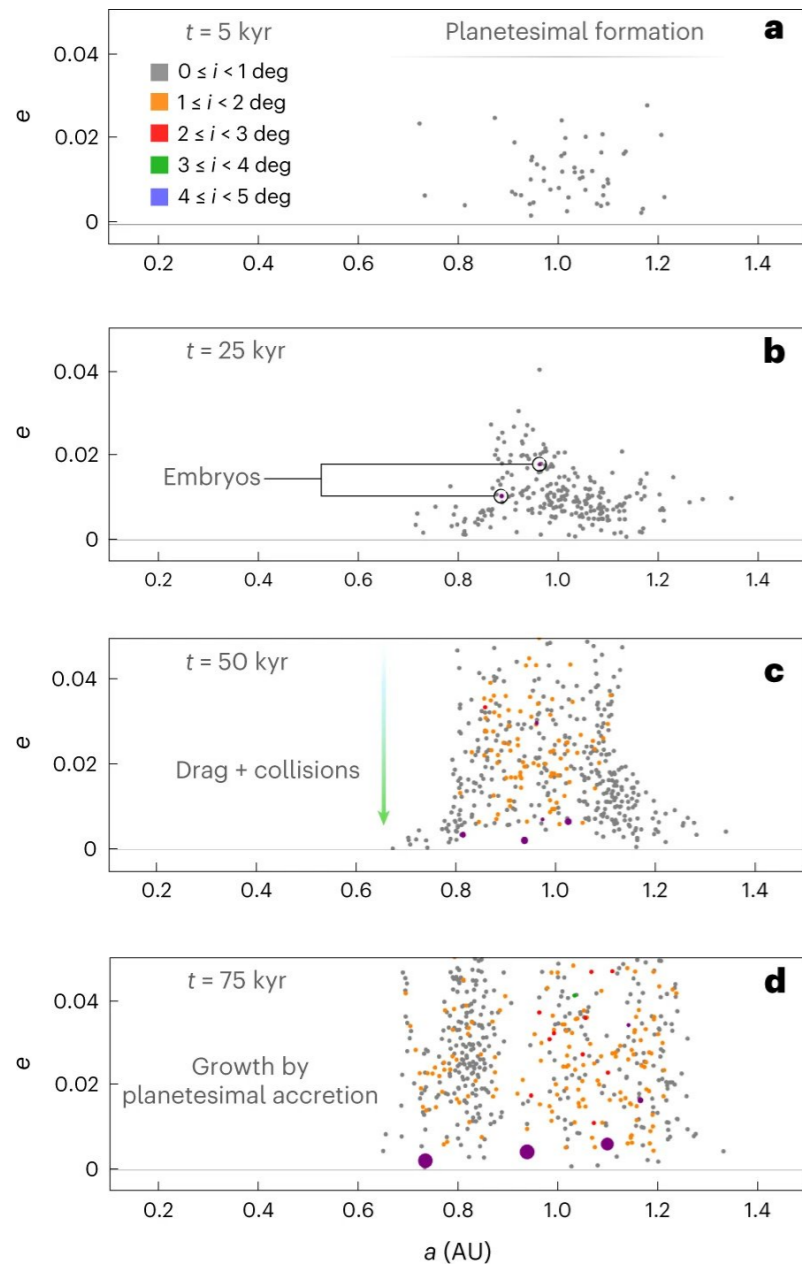




# Formation of rocky super-earths from a narrow ring of planetesimals

Konstantin Batygin  & Alessandro Morbidelli  *Nature Astronomy* **7**, 330–338 (2023) | [Cite this article](#)

 silicate-rich composition     intra-system uniformity     typical mass  $\sim$  few  $\times$  Earth     link to Jup, Sol



Orbital migration delivers planets to the disk's inner edge, where they stabilize. Thus, at face-value, the data appears to suggest that disks are truncated at an orbital period of  $\sim 3$  days, independent of the central mass...

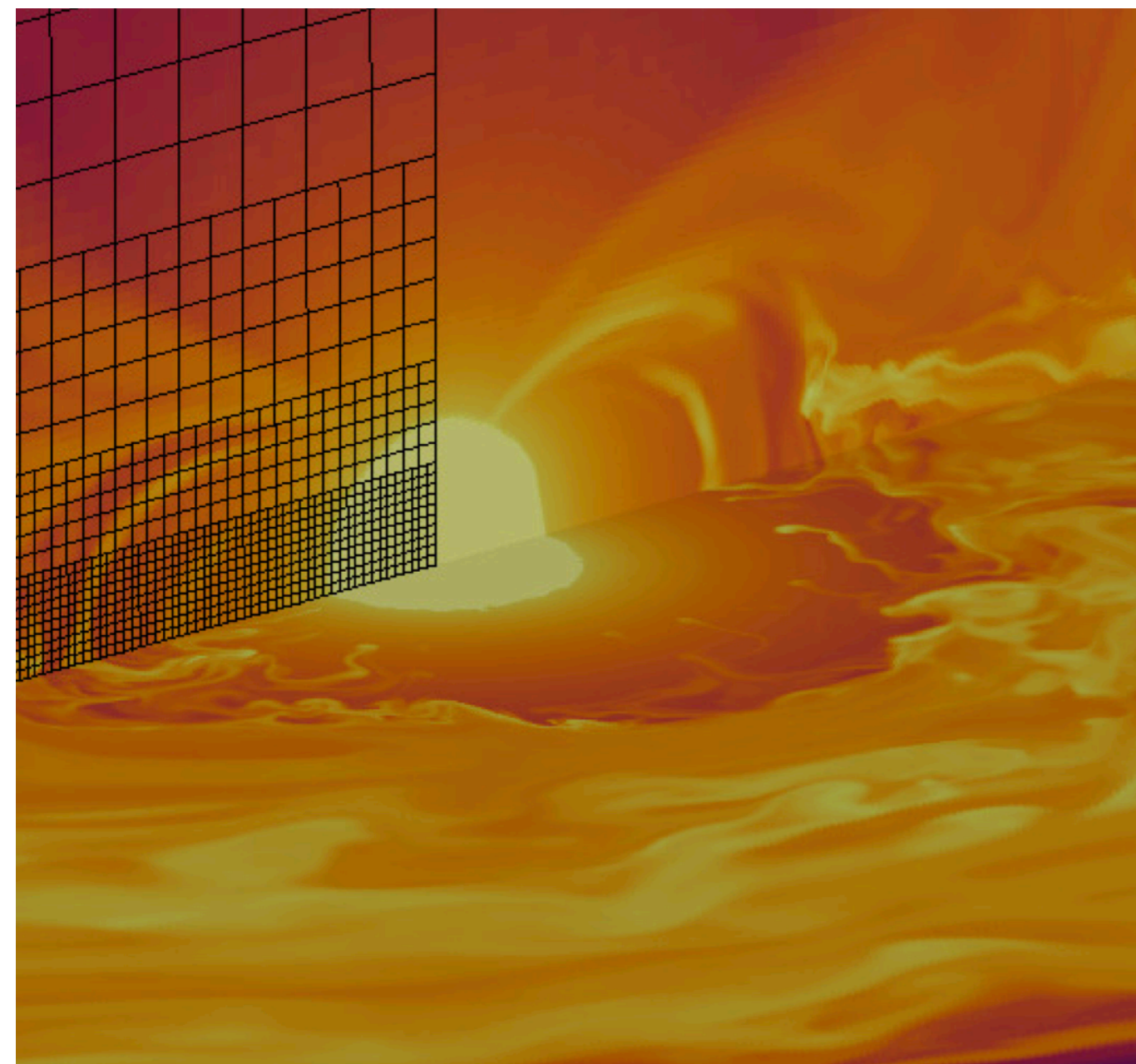


Fig. from Zhu et al (2023)

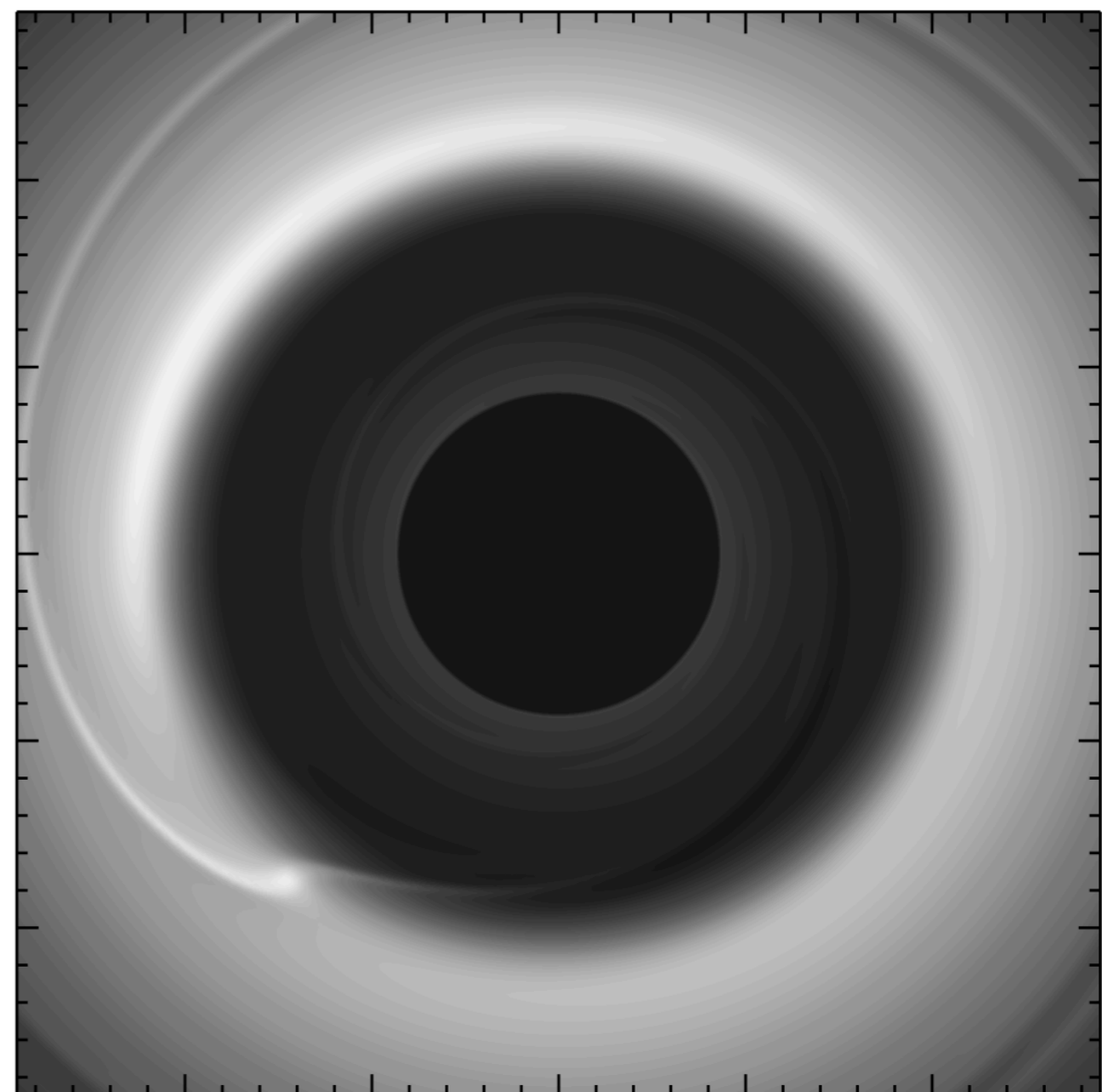
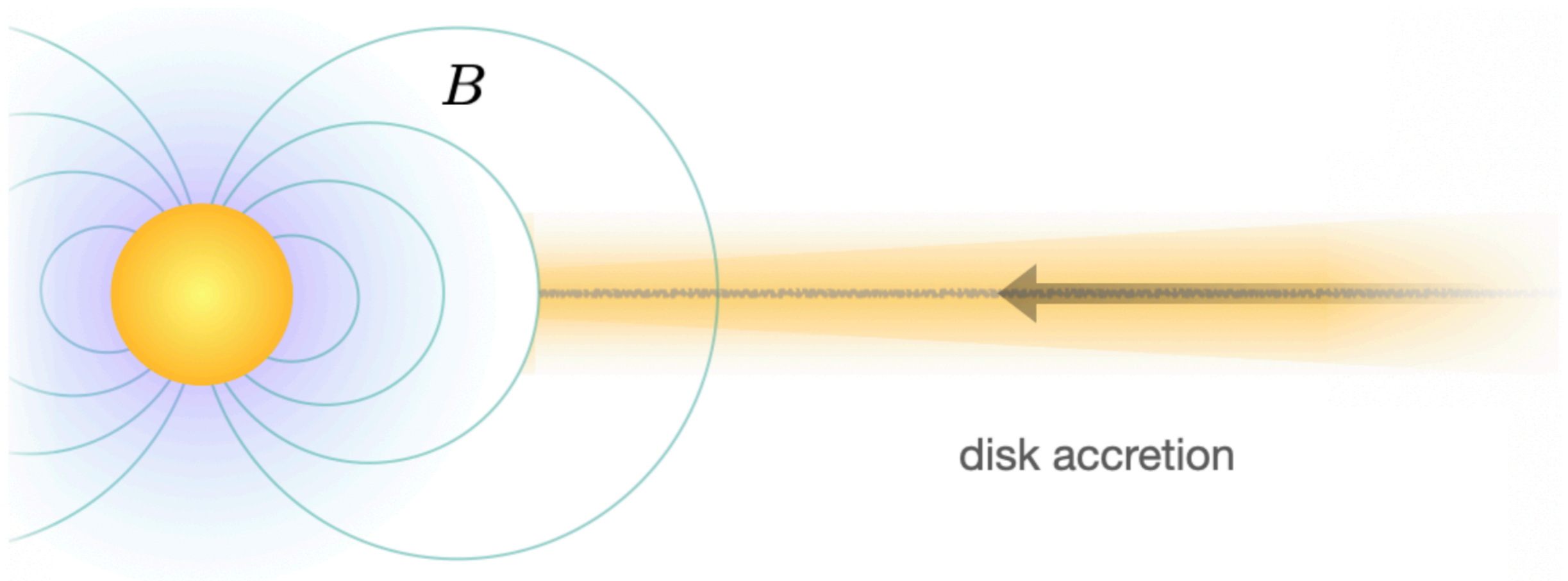


Fig. from Masset et al (2006)

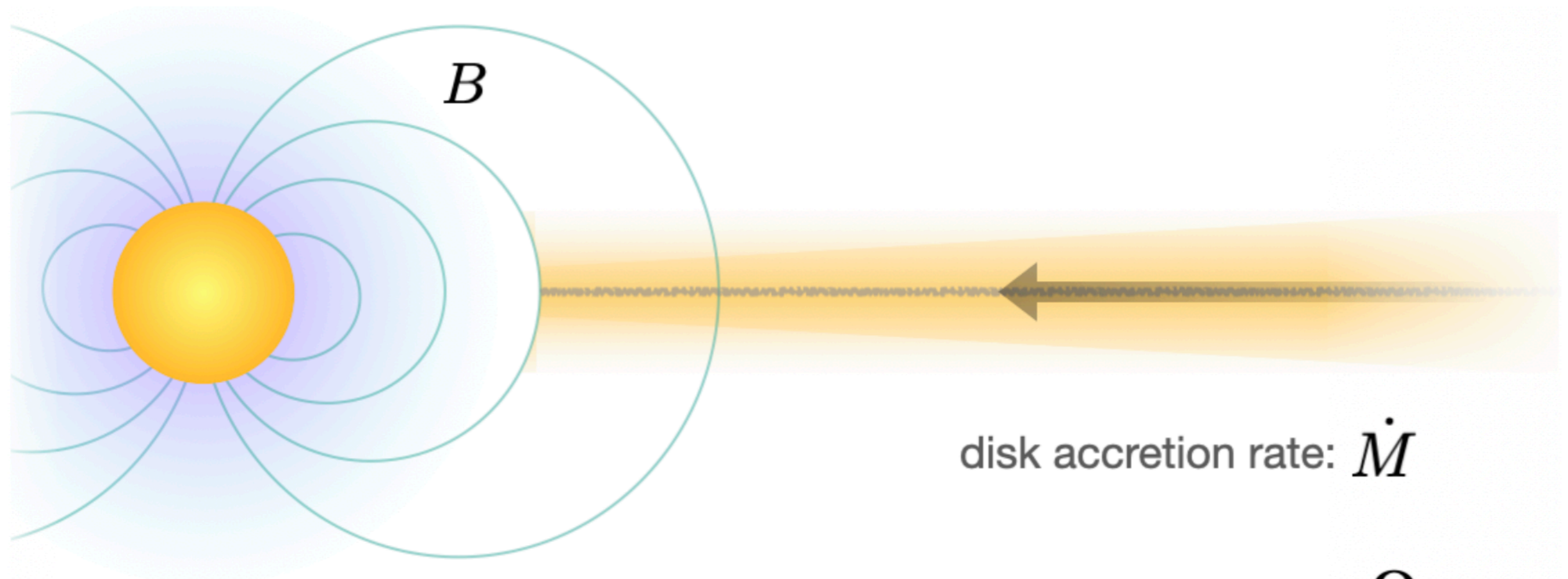


disk accretion

magnetospheric truncation

$$P_{\text{mag}} \sim P_{\text{ram}}$$

We need a theory that will connect stellar field ( $B$ ), radius ( $R$ ), accretion rate ( $\dot{M}$ ) etc.

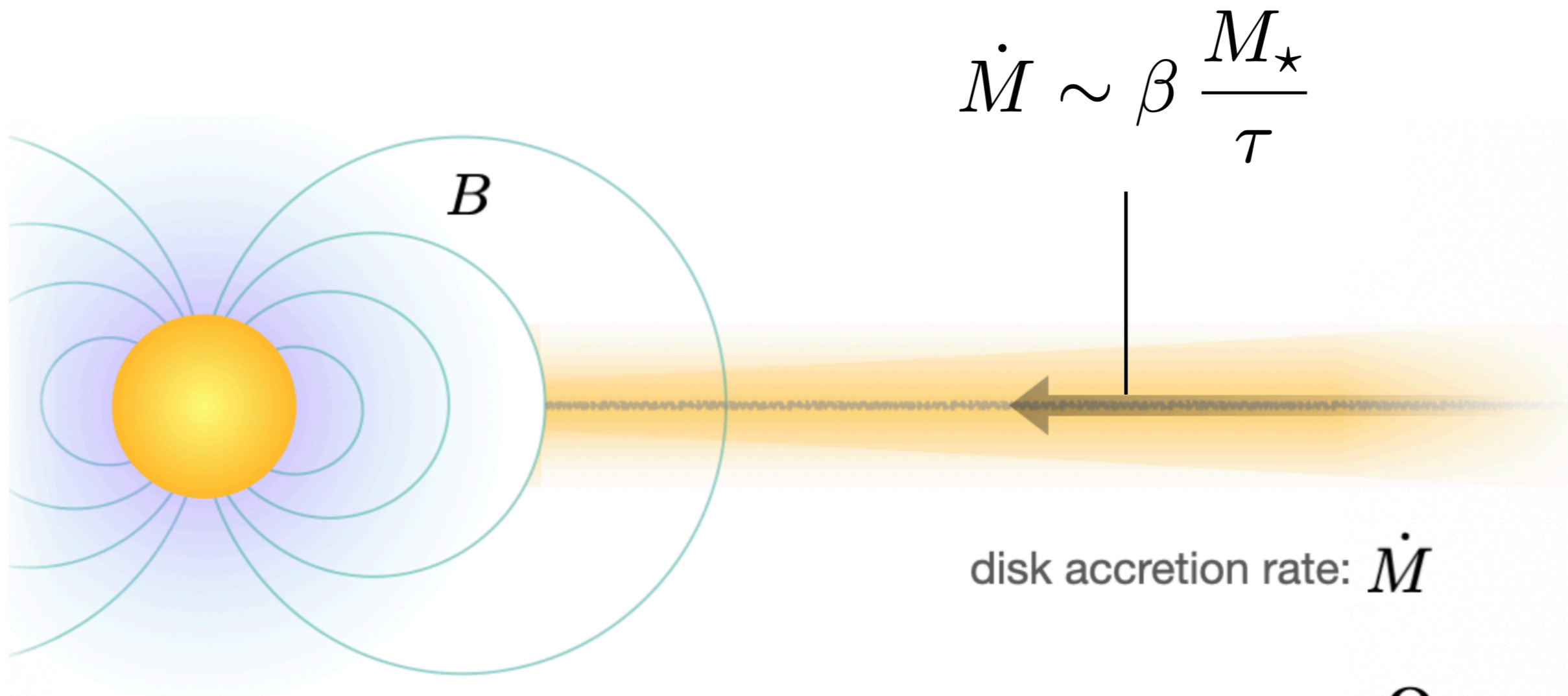


disk accretion rate:  $\dot{M}$

magnetospheric truncation frequency:  $\Omega$

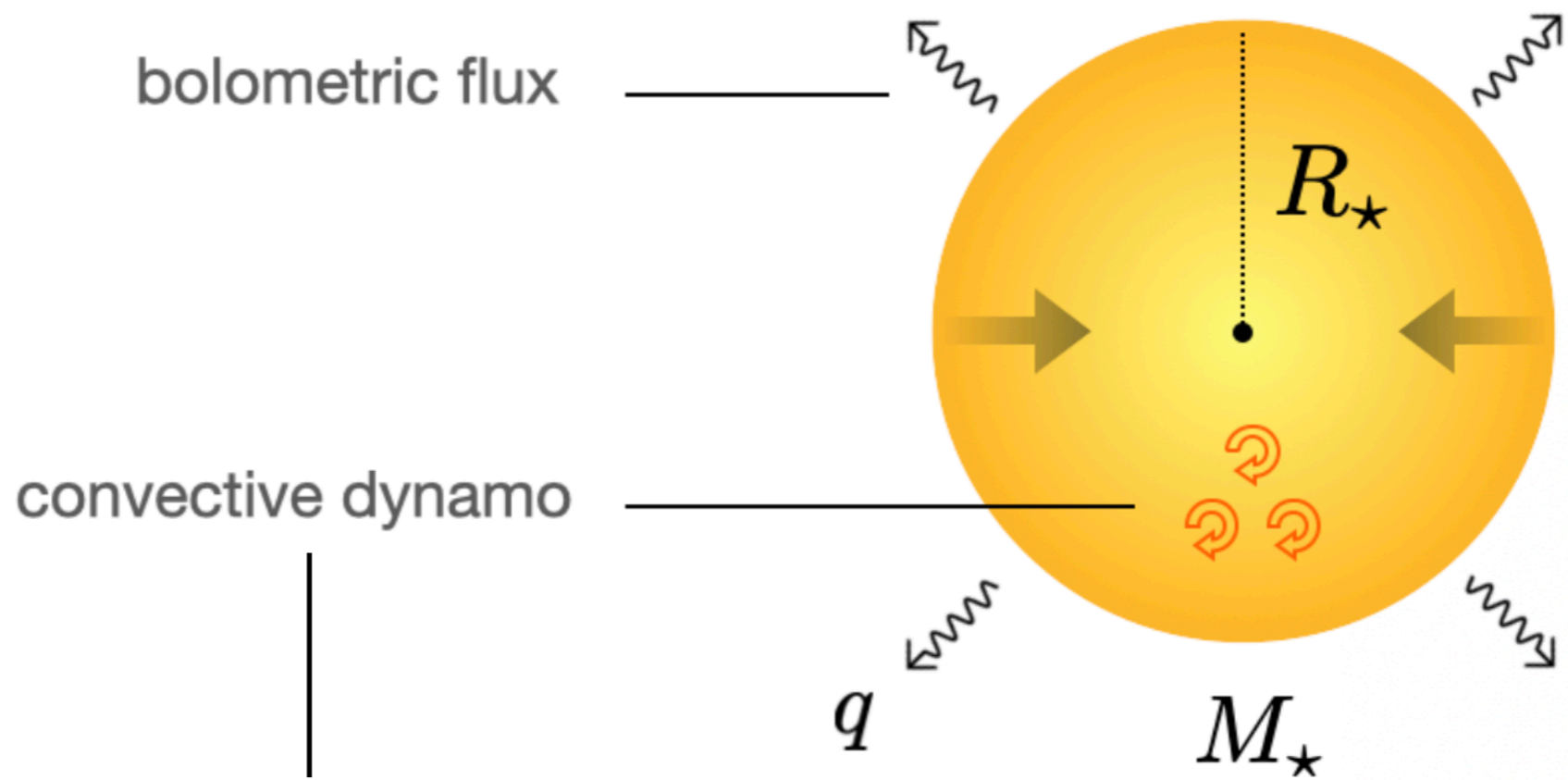
$$\frac{B^2}{2\mu_0} \sim \frac{B_\star^2}{2\mu_0} \left( \frac{R_\star}{r} \right)^6 \sim \frac{\dot{M}}{4\pi r^2} \sqrt{\frac{2\mathcal{G} M_\star}{r}}$$





magnetospheric truncation frequency:  $\Omega$

$$\frac{B_{\star}^2}{2\mu_0} \left( \frac{R_{\star}}{r} \right)^6 \sim \frac{\dot{M}}{4\pi r^2} \sqrt{\frac{2\mathcal{G} M_{\star}}{r}}$$



$$\frac{\langle B \rangle^2}{2\mu_0} \sim \rho v_{\text{conv}}^2 = c f_{\text{ohm}} \langle \rho \rangle^{1/3} (\mathcal{F} q)^{2/3}$$

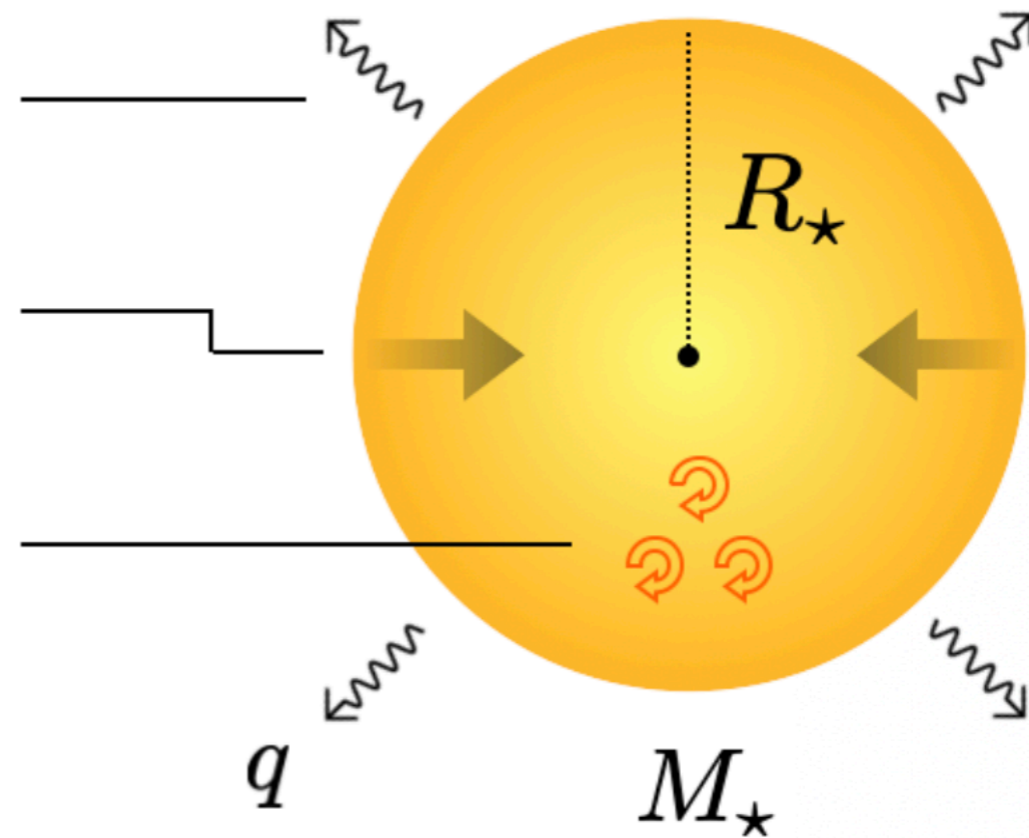
$$R_{\star} \approx \left( \frac{b G M_{\star}^2}{12 \pi q \tau} \right)^{1/3}$$

bolometric flux

central body undergoing gravitational  
(Kelvin-Helmholtz) contraction

convective dynamo

$$\frac{\langle B \rangle^2}{2 \mu_0} = c f_{\text{ohm}} \langle \rho \rangle^{1/3} (\mathcal{F} q)^{2/3}$$



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$$\Omega = 2 \xi \left[ \frac{\sqrt{2}}{(3 b \mathcal{F})^2} \left( \frac{\pi \beta \gamma^2}{c f_{\text{ohm}}} \right)^3 \frac{(\mathcal{G} \langle \rho \rangle)^3}{\tau} \right]^{1/7}$$

$$\approx 2.4 \times 10^{-5} \text{ s}^{-1} \approx \frac{2 \pi}{3 \text{ day}}$$